

Mechanism of Electronically Enhanced Defect Reactions in Semiconductors

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Keywords: defect reaction, lattice relaxation, multiphonon process, phonon kick, coherent capture

Electronic transitions sometimes induce or enhance defect reactions in semiconductors, such as impurity diffusion, structural change of a defect (creation, annihilation and multiplication), and climb and glide motions of dislocation [1]. It has been suggested that the transient vibration induced by a carrier capture enhances the probability of a following capture of the opposite carrier [2,3]. Recently Shinozuka and Karatsu have simulated the transient lattice vibration induced by successive captures and its effect on the next capture process [4]. If N pairs of electrons and holes are sequentially captured by a defect, each occurred at $t_{e1} < t_{h1} < t_{e2} < t_{h2} < \dots < t_{eN} < t_{hN}$, the time evolution of $Q_1(t)$ for $t_{hN} < t$ is given by

$$Q_1(t) = Q_1^{\text{before}}(t) + \bar{Q}_1 \left\{ - \sum_j^N \exp[-\Delta\omega^2(t-t_{ej})^2/4] \cos\omega_0(t-t_{ej}) + \sum_j^N \exp[-\Delta\omega^2(t-t_{hj})^2/4] \cos\omega_0(t-t_{hj}) \right\}, \quad (1)$$

for a system whose phonon frequency ω_k distribution is given by a Gaussian with the central frequency ω_0 and the width $\Delta\omega$. Here $Q_1^{\text{before}}(t)$ is the time dependence if all the captures of N electron-hole pairs would not have occurred at all. Thus, $Q_1(t)$ is shown to be a linear combination of damping oscillations, each starts at time t_{ej} or t_{hj} and lasts in the same duration $\tau \sim 2\pi/\Delta\omega$. It should be reminded that the capture time t_{ej} and t_{hj} do not take an arbitrary value but should be obeyed by probability.

We will discuss the effect of the transient vibration (eq. 1) to defect reactions using following assumptions:

- The first electron capture takes place by a thermal activation Q_1 to Q_e with the activation energy E_e^{act} .
- If $Q_1(t)$ crosses at Q_e (Q_h), there is a probability p_e (p_h) per time to capture an electron (hole) nonradiatively. p_e (p_h) $\sim \omega_0/2\pi$ in the adiabatic limit.
- The defect reaction coordinate is given by $Q_R(t) = \sum_k f_k q_k(t)$, whose frequency distribution is a Gaussian with the central frequency ω_R and the width $\Delta\omega$.
- If $Q_R(t)$ crosses the critical point Q_c , the defect reaction takes place.

We have found that every capture process enhances the vibration of $Q_1(t)$. On the other hand, it is not so for $Q_R(t)$ in $\omega_0 \neq \omega_R$ case, because the timing of push and pull for $Q_R(t)$ is out of phase (Fig.1).

Let us discuss the possibility of recombination-

enhanced defect reaction in connection with the transient induced lattice vibration. The condition for the coherent captures turns out to be $p_e \tau \sim p_h \tau \sim \omega_0 \Delta\omega > 1$. Then if N pairs of electron and hole are captured within a short period $\sim 2\pi/\Delta\omega$, the amplitude of the interaction mode $Q_1(t)$ increases remarkably. If the central frequency ω_0 of the reaction coordinate Q_R is not so much different from ω_0 of the interaction mode Q_1 , more than the band gap energy $E_g = E_e^{\text{th}} + E_h^{\text{th}}$ can be transformed by a series of coherent carrier captures into the lattice vibration energy. The defect reaction rate is given by $(\omega_0/2\pi)\exp(-E_i^{\text{act}}/k_B T)$ because only the first capture ($i=e, h$) is to be activated. On the other hand, if ω_R is much different from ω_0 the rate is $(\omega_0/2\pi)\exp(-U_0^*/k_B T)$ with $U_0^* = (E_i^{\text{act}} + E_i^{\text{th}})$ because the N phonon-kick's are out of phase.

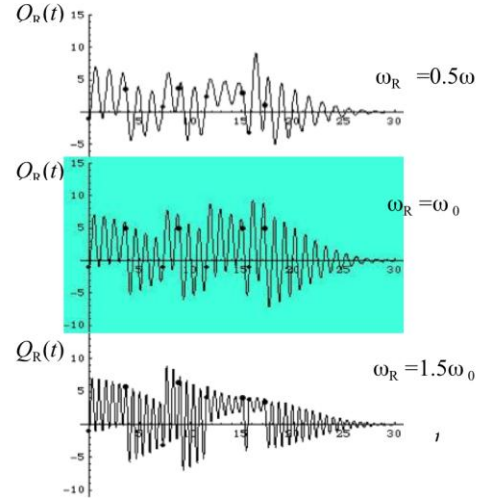


Fig. 1 The time evolution of $Q_R(t)$. Large (small) circle indicates the capture time t_{ei} (t_{hi}) for an electron (hole).

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