

# Modelling of bulk superconductor magnetisation: a review

Dr Mark Ainslie

*Engineering & Physical Sciences Research Council (EPSRC)  
Research Fellow*

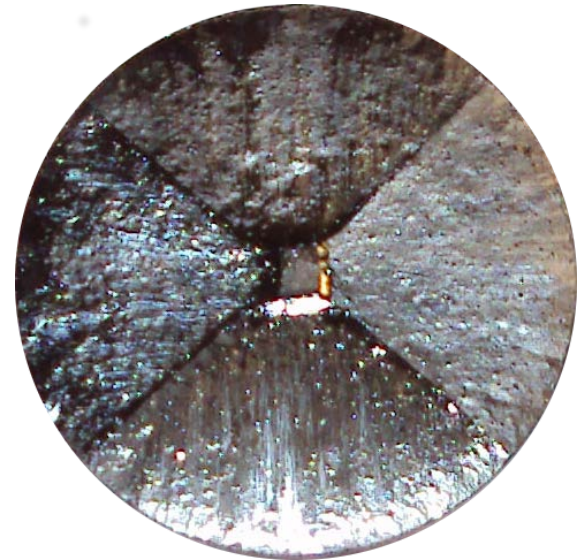
Bulk Superconductivity Group, Department of Engineering

# Presentation Outline

- **Overview of modelling bulk superconductor magnetisation**
- **Case studies:**
  - Multi-pulse, pulsed field magnetisation (PFM) of bulk high-temperature superconductors
  - Split coil PFM with an iron yoke
  - Field cooling (FC) magnetisation of iron-pnictide (Ba122) bulks

# Bulk Superconductors

- Bulk superconducting materials can 'trap' large magnetic fields  $> 17\text{ T}$
- Achieved by pinning penetrated magnetic field (quantised flux lines)  $\rightarrow$  macroscopic electrical currents
- Magnetisation increases with sample volume



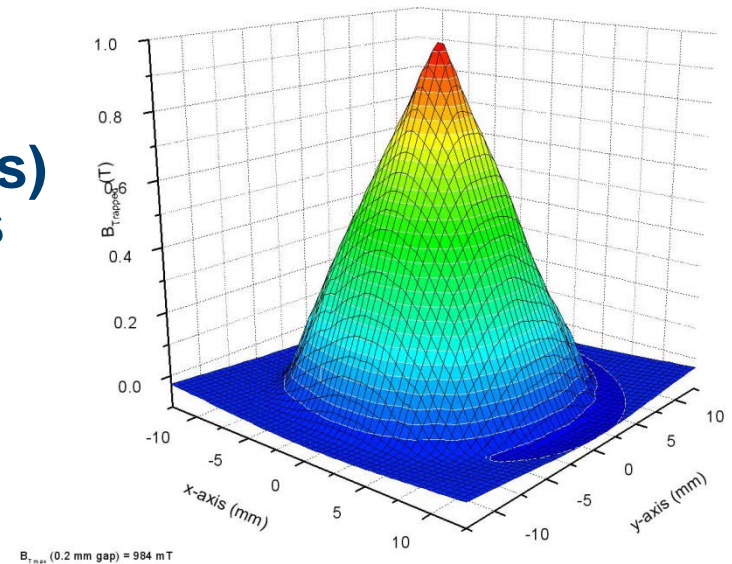
A large, single grain bulk superconductor

# Bulk Superconductors

- Bulk superconducting materials can 'trap' large magnetic fields > 17 T
- Achieved by pinning penetrated magnetic field (quantised flux lines) → macroscopic electrical currents
- Magnetisation increases with sample volume
- Trapped field given by

$$B_{\text{trap}} = k \mu_0 J_c R$$

$$\text{where } k = \frac{t_B}{2R} \cdot \ln \left( \frac{R + \sqrt{R^2 + t_B^2}}{t_B} \right)$$



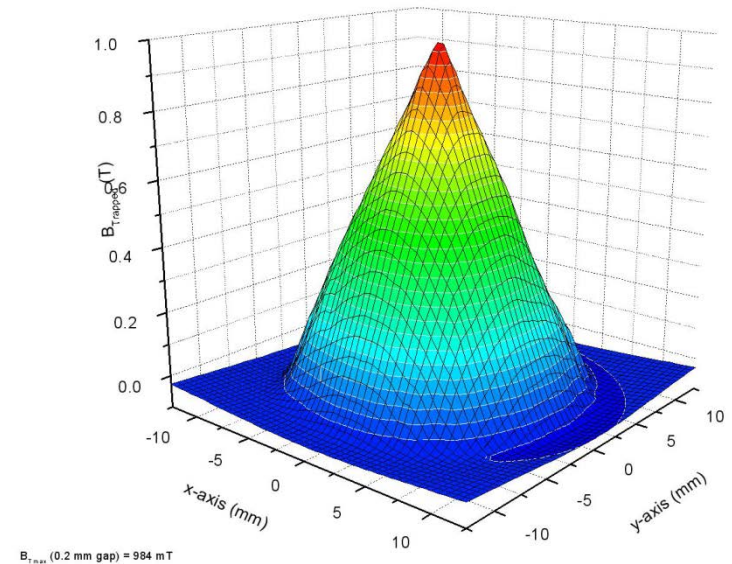
Typical trapped magnetic field profile of a bulk superconductor

# Bulk Superconductors

$$B_{\text{trap}} = k \mu_0 J_c R$$

$$\text{where } k = \frac{t_B}{2R} \cdot \ln \left( \frac{R + \sqrt{R^2 + t_B^2}}{t_B} \right)$$

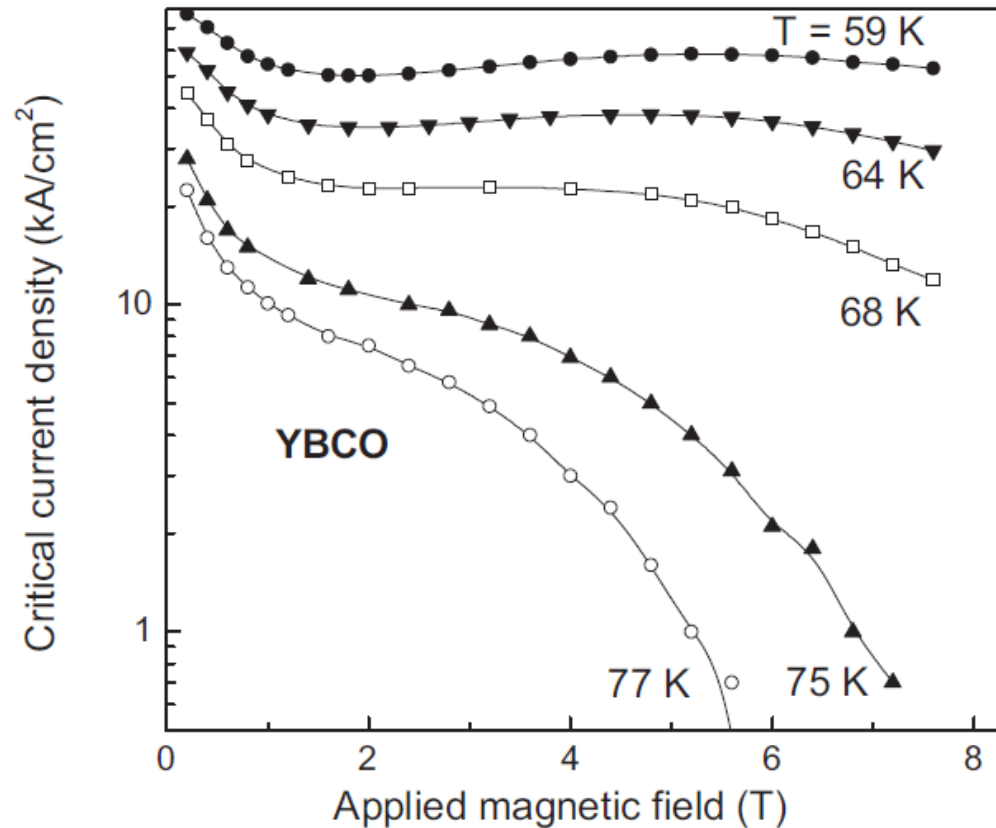
- From Bean model (infinite slab) + Biot-Savart law
- Example of an analytical model
  - Easier to deal with & faster
  - Rely on specific simplified geometries & simplified, homogeneous assumptions:
    - Constant / uniform  $J_c$ , no frequency / time dependence, etc.



# Bulk Superconductors

$$B_{\text{trap}} = k \mu_0 J_c R$$

- **Candidate materials must be able to:**
  - Pin magnetic flux effectively
  - Carry large current density,  $J_c$ , over large length scales
  - Be insensitive to application of large magnetic fields,  $J_c(B)$



**Example field dependence of critical current density,  $J_c(B)$ , for bulk YBCO**

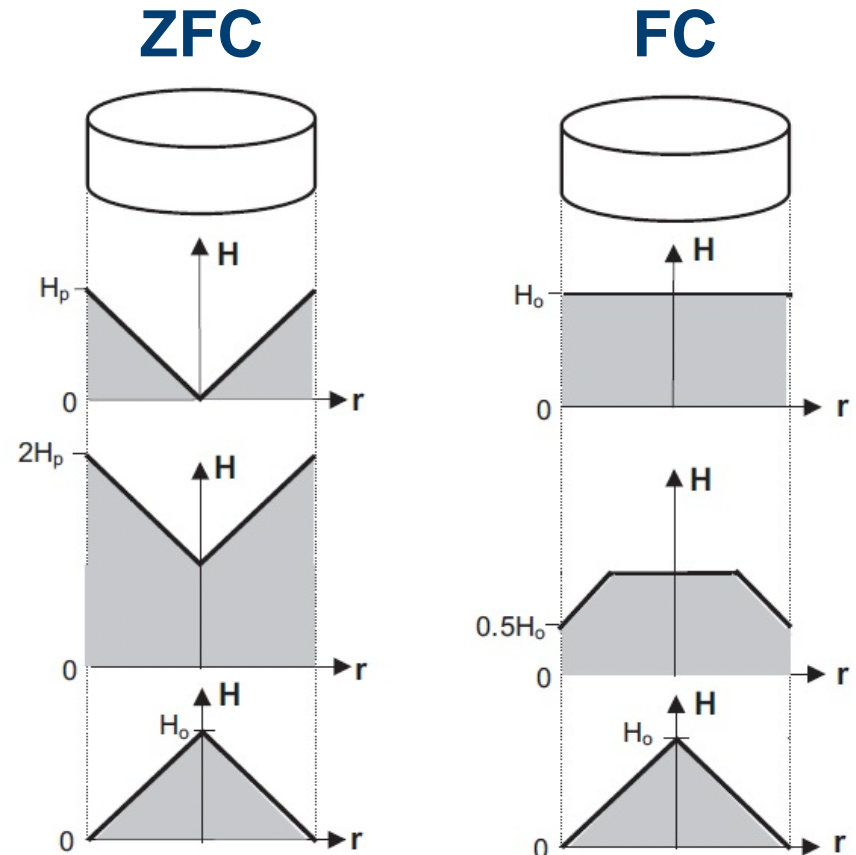
# Magnetisation of Bulk Superconductors

- Three magnetisation techniques:

- Field Cooling (FC)
- Zero Field Cooling (ZFC)
- Pulsed Field Magnetisation (PFM)

- To trap  $B_{\text{trap}}$ , need at least  $B_{\text{trap}}$  or higher

- FC and ZFC require large magnetising coils, long magnetizing times
- Impractical for applications/devices → PFM



# Numerical Modelling of Magnetisation

- **Numerical models can simulate practical & complex situations, playing a number of crucial roles:**
  - Simulate accurate magnetic field, current, temperature, mechanical stress distributions
  - Interpret experimental results & physical mechanisms of bulk superconductor magnetisation
  - Design & predict performance of magnetising fixtures & techniques
  - Design & predict performance of practical bulk superconductor-based devices



# Numerical Modelling of Magnetisation

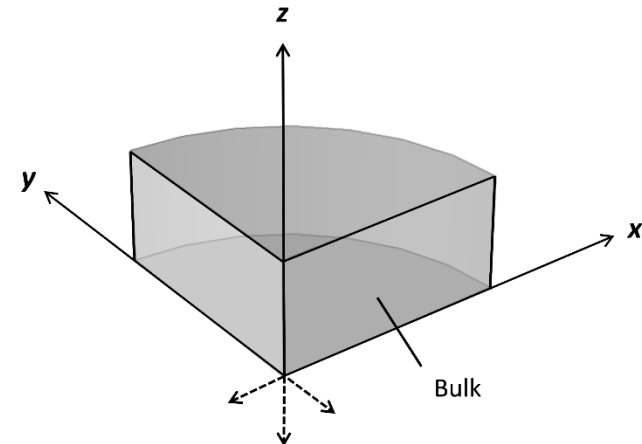
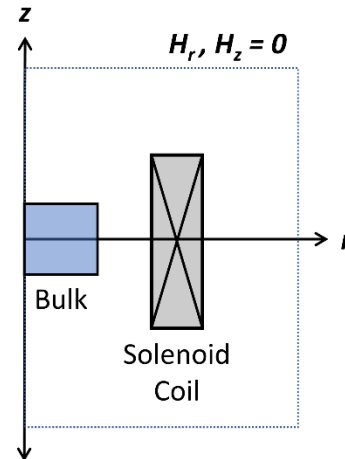
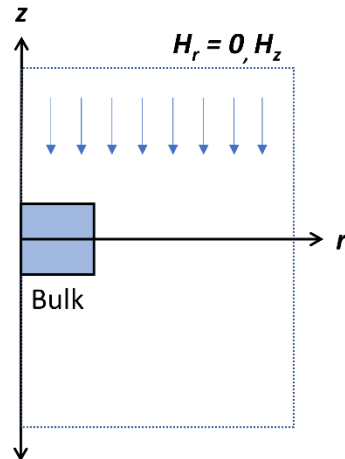
BULK GEOMETRY &  
MAGNETISATION  
FIXTURE

ELECTROMAGNETIC  
FORMULATION

THERMAL  
EQUATIONS &  
PROPERTIES

$J_c(B, T)$

$E$ - $J$  POWER LAW



- 2D axisymmetric generally sufficient for cylindrical bulks with a homogeneous  $J_c$  distribution
- 3D required for an inhomogeneous  $J_c$  distribution around the  $ab$ -plane; for non-symmetric shapes

# Numerical Modelling of Magnetisation

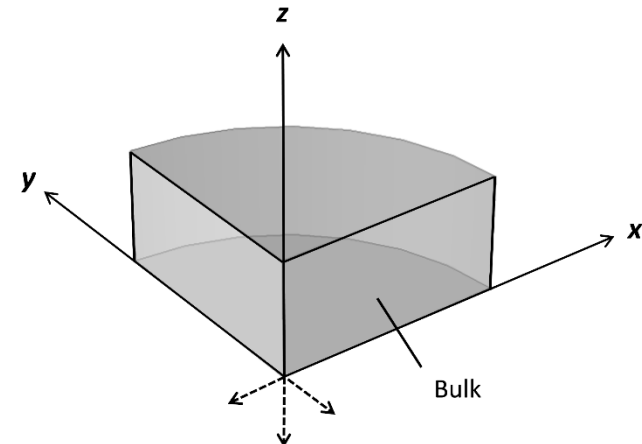
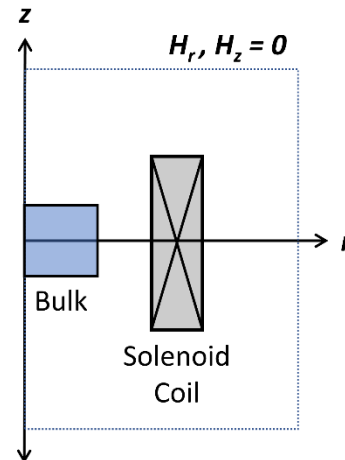
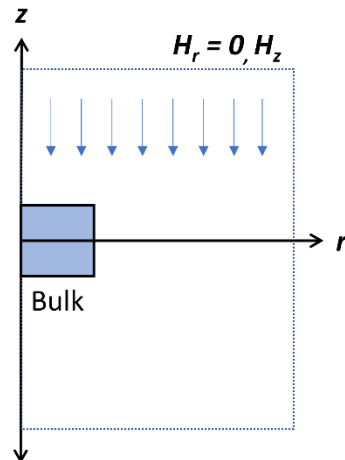
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- Magnetising fixture: uniform boundary conditions or inserting a copper coil sub-domain
- Cooling: using a cold head + vacuum chamber or submersion in liquid cryogen

# Numerical Modelling of Magnetisation

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Finite element method is commonly used & well developed (other techniques do exist)

Governing equations:

Maxwell's equations ( $\mathbf{H}$  formulation)

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} = -\frac{\partial (\mu_0 \mu_r \mathbf{H})}{\partial t} \quad \text{Faraday's Law}$$

$$\nabla \times \mathbf{H} = \mathbf{J} \quad \text{Ampere's Law}$$

Other formulations also exist ( $\mathbf{A}$ - $V$ ,  $T$ - $\Omega$ , *Campbell's equation*)

# Numerical Modelling of Magnetisation

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$E$ - $J$  POWER LAW

**Thermal behaviour needs to be modelled when the bulk experiences a significant change in temperature**

e.g., during PFM, modelling complete FC magnetisation process

**Governing equations:**

$$\rho \cdot C \frac{dT}{dt} = \nabla \cdot (k \nabla T) + Q$$

$$Q = E \cdot J$$

$\rho$  = mass density,  $C$  = specific heat,  $\kappa$  = thermal conductivity,  $Q$  = heat source

# Numerical Modelling of Magnetisation

BULK GEOMETRY &  
MAGNETISATION  
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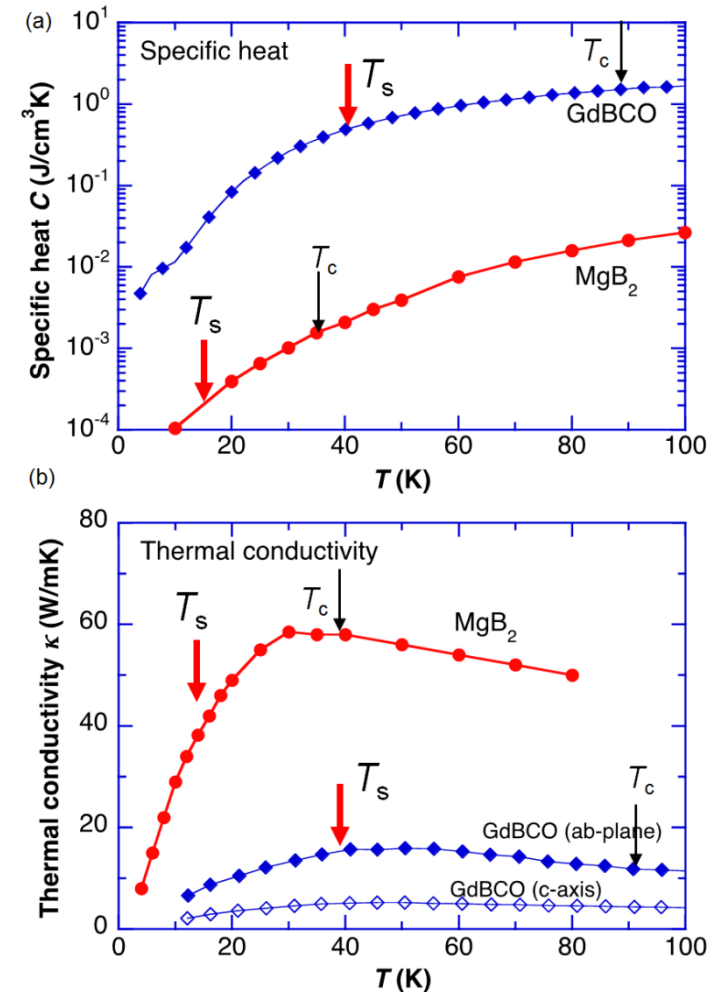
THERMAL  
EQUATIONS &  
PROPERTIES

$J_c(B, T)$

$E$ - $J$  POWER LAW

Can use measured  
experimental data + fitting  
function or interpolation  
over a specific temperature  
range

Can choose constant  
parameters for  $C$ ,  $\kappa$  for  
 $T = T_{op}$  as a reasonable  
approximation



# Numerical Modelling of Magnetisation

BULK GEOMETRY &  
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$J_c(B, T)$

$E$ - $J$  POWER LAW

## HTS materials Kim-like model:

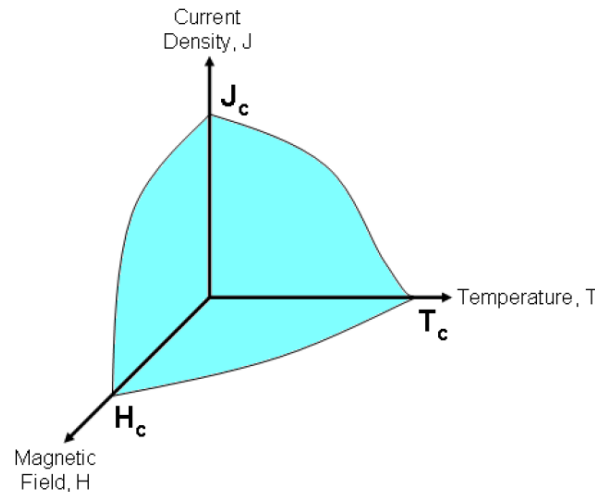
$$J_c = \frac{J_{c0}}{\left(1 + \frac{B}{B_0}\right)^\alpha}$$

**Fishtail effect:**  $J_c(B) = J_{c1} \exp\left(-\frac{B}{B_L}\right) + J_{c2} \frac{B}{B_{\max}} \exp\left[\frac{1}{y} \left(1 - \left(\frac{B}{B_{\max}}\right)^y\right)\right]$

## MgB<sub>2</sub> materials

$$J_c(B, T) = J_{c0}(T) \exp\left(-\frac{B}{B_0}\right)^a$$

$$J_{c0}(T) = \alpha \left[1 - \left(\frac{T}{T_c}\right)^2\right]^{1.5}$$



# Numerical Modelling of Magnetisation

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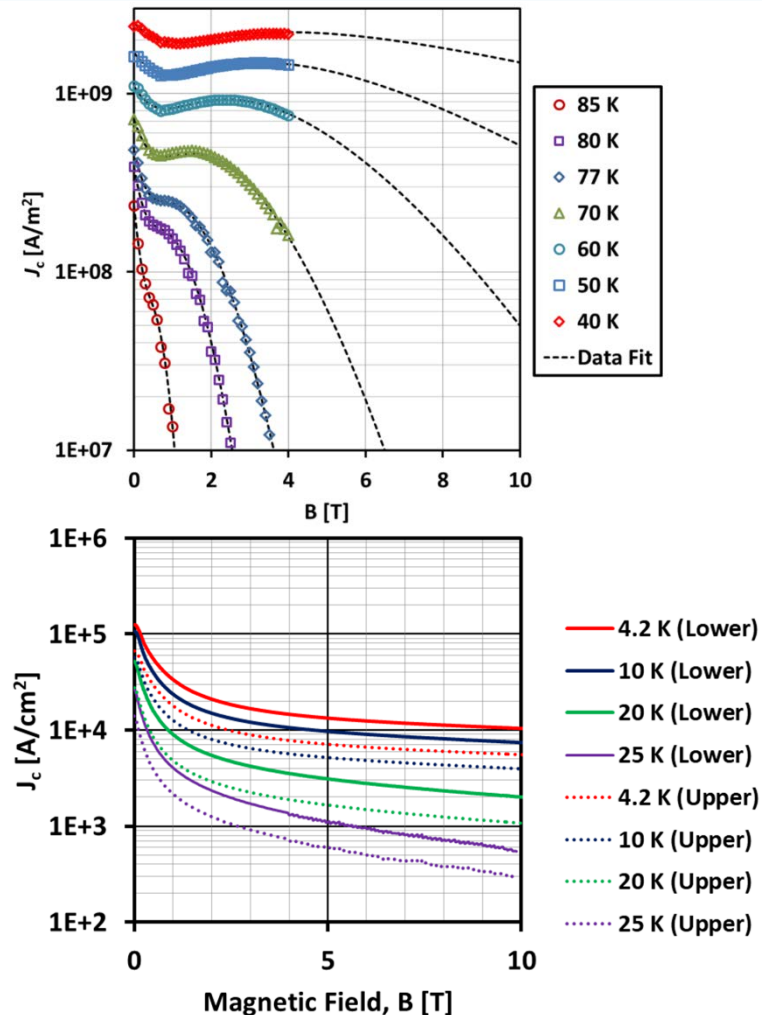
$J_c(B, T)$

$E$ - $J$  POWER LAW

... or direct interpolation

$\text{GdBa}_2\text{Cu}_3\text{O}_{7-\delta}$  (15wt% Ag)

$\text{Ba}_{0.6}\text{K}_{0.4}\text{Fe}_2\text{As}_2$  (Ba122)



# Numerical Modelling of Magnetisation

BULK GEOMETRY &  
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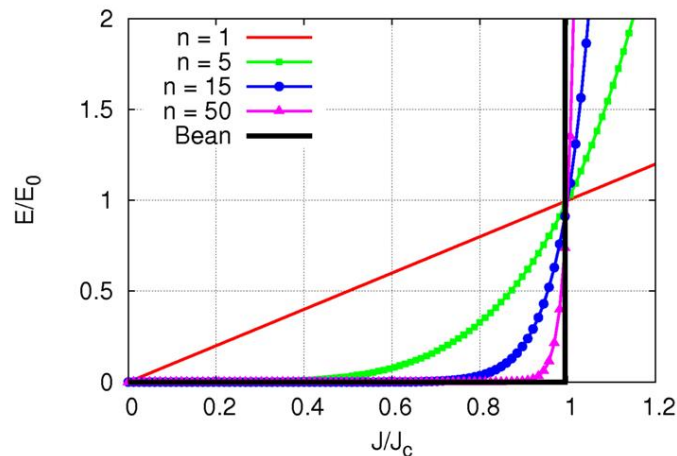
THERMAL  
EQUATIONS &  
PROPERTIES

$J_c(B, T)$

***E-J* POWER LAW**

## ***E-J* power law**

- Conventional materials → non-linear permeability, linear resistivity
- Superconductors → linear permeability ( $\mu_0$ ), non-linear resistivity
- Non-linearity is extreme: power law with  $n > 20$



***I-V* curves for different  $n$  values**

$$\mathbf{E} = E_0 \left( \frac{J}{J_c} \right)^{n-1} \frac{J}{J_c}$$

$$E = \rho J$$

$$\mathbf{E} = E_0 \left( \frac{J}{J_c} \right)^{n-1} \frac{J}{J_c}$$



## Case study #1:

*Multi-pulse, pulsed field magnetisation  
(PFM) of bulk high-temperature  
superconductors*

# Pulsed Field Magnetisation

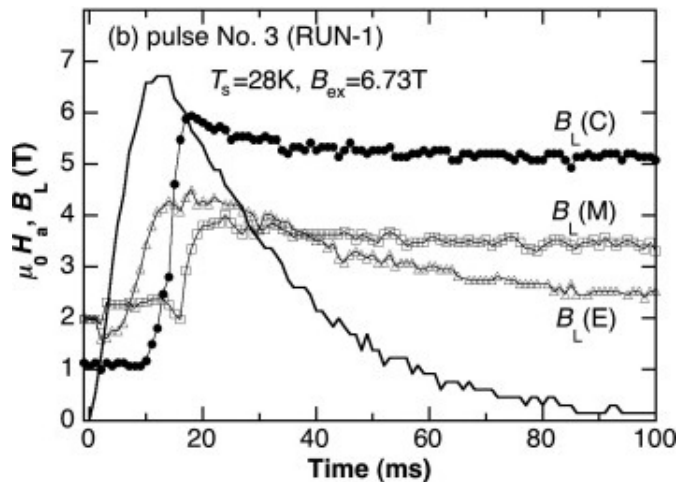
- **PFM technique: compact, mobile, relatively inexpensive**
- **Main issue:  $B_{\text{trap}}$  [PFM] <  $B_{\text{trap}}$  [FC], [ZFC]**
  - Temperature rise  $\Delta T$  due to rapid movement of magnetic flux
- **Record PFM trapped field: 5.2 T @ 29 K**

Top surface of 45 mm diameter Gd-Ba-Cu-O  
Fujishiro et al. *Physica C* 2006
- **Record trapped field by FC: 17.6 T @ 26 K**

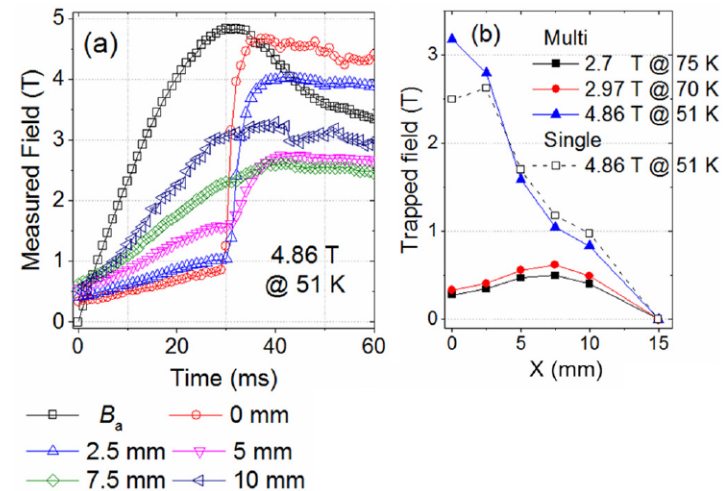
Centre of 2 x 25 mm diameter Gd-Ba-Cu-O  
Durrell et al. *Supercond. Sci. Technol.* 2014

# Pulsed Field Magnetisation

- Many considerations for PFM:
  - Pulse magnitude, pulse duration, temperature(s), number of pulses, type of magnetising coil(s), use of ferromagnetic materials
  - Dynamics of magnetic flux during PFM process
- Multi-pulse PFM: effective in increasing trapped field/flux



Fujishiro et al. *Physica C* 2006

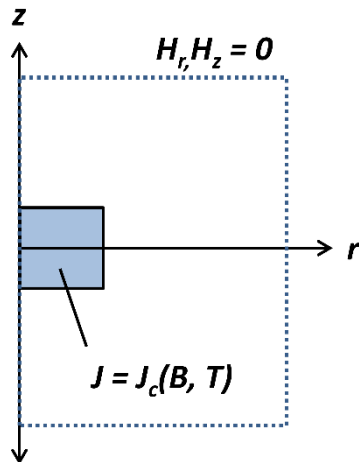


Zhou et al. *Appl. Phys. Lett.* 2017

# Modelling Trapped Field Capability (FC/ZFC)

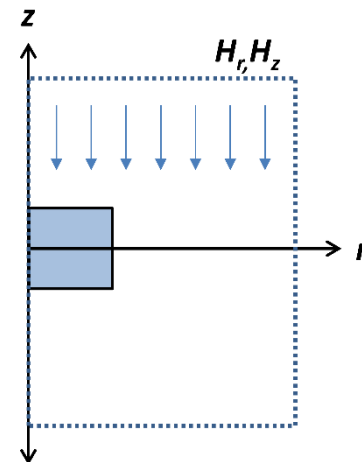
## Model #1

- Stationary
- Comsol Multiphysics
  - 2D axisymmetric
  - AC/DC module
  - Magnetic Field (mf) interface
  - External Current Density node
- No flux creep
- Time taken: ~ 2-3 seconds



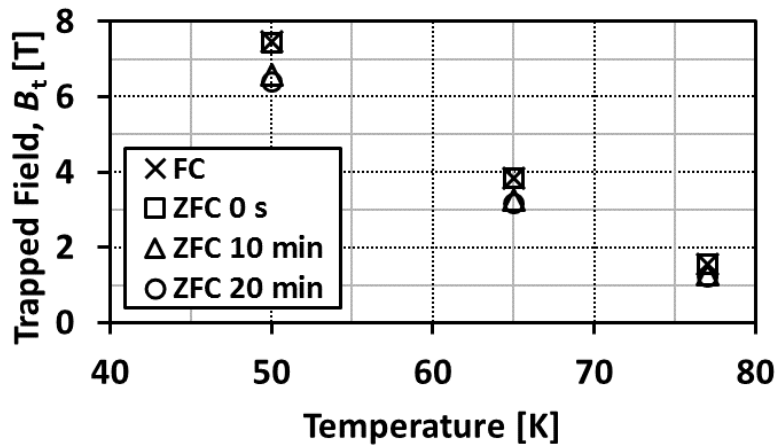
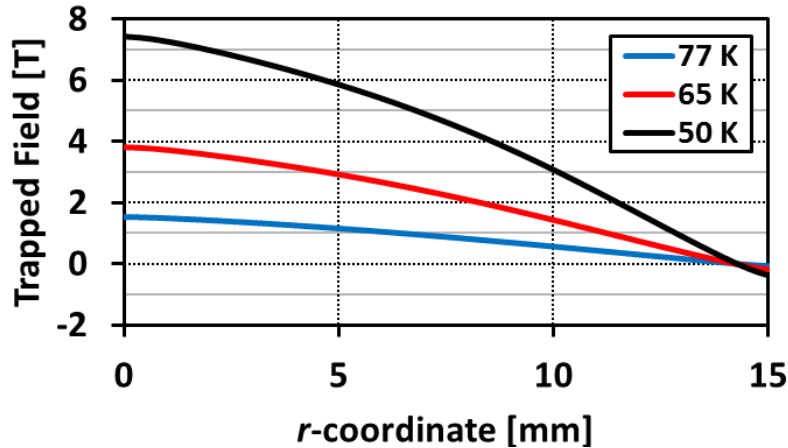
## Model #2

- Time-dependent
- Comsol Multiphysics
  - 2D axisymmetric
  - AC/DC module
  - Magnetic Field Formulation (mfh) interface
- $E$ - $J$  power law,  $E \propto J^n$  (flux creep)
- Apply + remove background field
- Time taken: up to 1-2 hours



# Modelling Trapped Field Capability

## Results



GdBa<sub>2</sub>Cu<sub>3</sub>O<sub>7-δ</sub> (15wt% Ag)

Magnetisation	Time	Trapped Field	
77 K			
FC	---	1.544 T	79%
ZFC [5 T]	t = 0 min	1.546 T	
	t = 10 min	1.263 T	
	t = 20 min	1.223 T	
65 K			
FC	---	3.826 T	82.5%
ZFC [10 T]	t = 0 min	3.827 T	
	t = 10 min	3.256 T	
	t = 20 min	3.158 T	
50 K			
FC	---	7.449 T	86%
ZFC [20 T]	t = 0 min	7.422 T	
	t = 10 min	6.577 T	
	t = 20 min	6.405 T	

# PFM Modelling Framework

Electromagnetic properties modelled as Model #2 (ZFC);  
magnetising fixture assumed as solenoid coil:

$$I_{pulse}(t) = N \cdot I_0 \frac{t}{\tau} \exp\left(1 - \frac{t}{\tau}\right) \quad \tau = 15 \text{ ms}$$

Thermal behaviour:

$$\rho \cdot C \frac{dT}{dt} = \nabla \cdot (k \nabla T) + Q$$

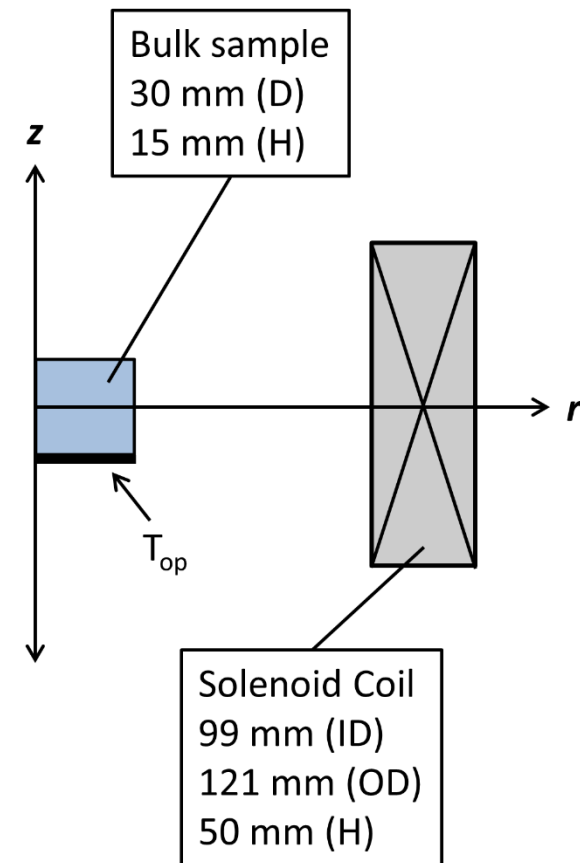
$$Q = E \cdot J \quad \text{Heat source, coupling with EM model + } J_c(B, T)$$

$\rho$  = mass density (bulk 5900 kg/m<sup>3</sup>, indium 7310 kg/m<sup>3</sup>)

$C$  = specific heat (measured, temperature-dependent)

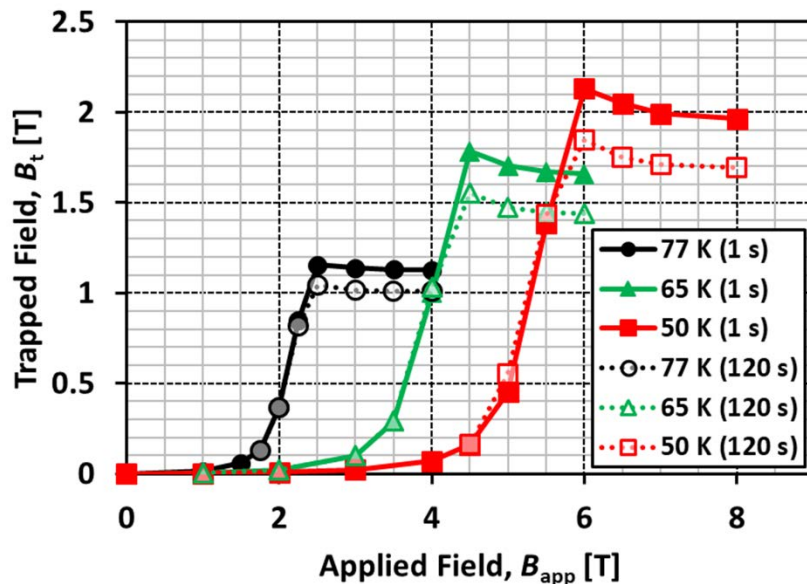
$\kappa$  = thermal conductivity:

$\kappa_{ab} = 20 \text{ W/(m}\cdot\text{K)}$ ,  $\kappa_c = 4 \text{ W/(m}\cdot\text{K)}$ ,  $\kappa_{\text{indium}} = 0.5 \text{ W/(m}\cdot\text{K)}$

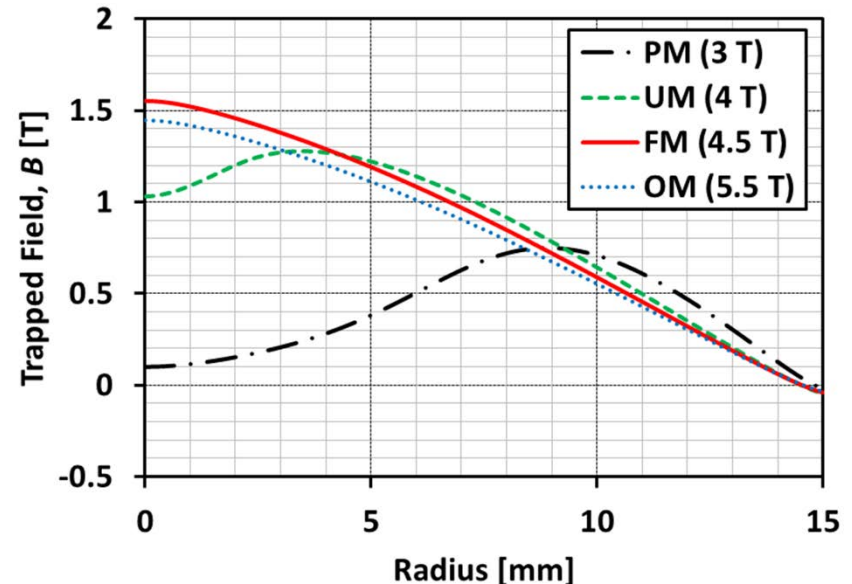


# PFM Single Pulse Results

- $t = 120$  s  $\rightarrow$  flux creep relaxation & cooling back to operating temp.
- Percentage of ZFC( $t = 20$  min):
  - 77 K 85%
  - 65 K 49%
  - 50 K 29%

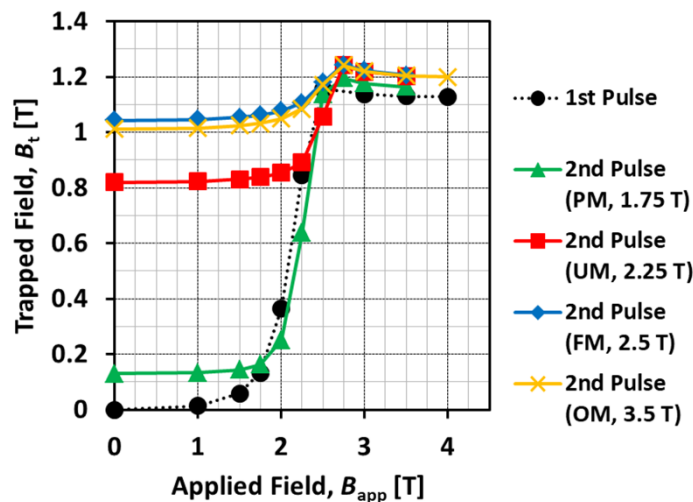


- Four specific cases as initial conditions for 2<sup>nd</sup> pulse:
  - Partially-magnetised (PM), so-called 'M-shaped' profile
  - Under-magnetised (UM)
  - Fully-magnetised (FM)
  - Over-magnetised (OM)

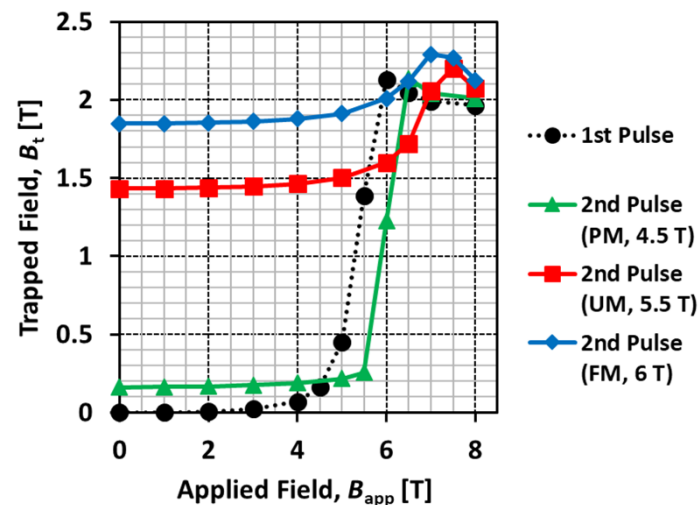


# PFM 2<sup>nd</sup> Pulse Results

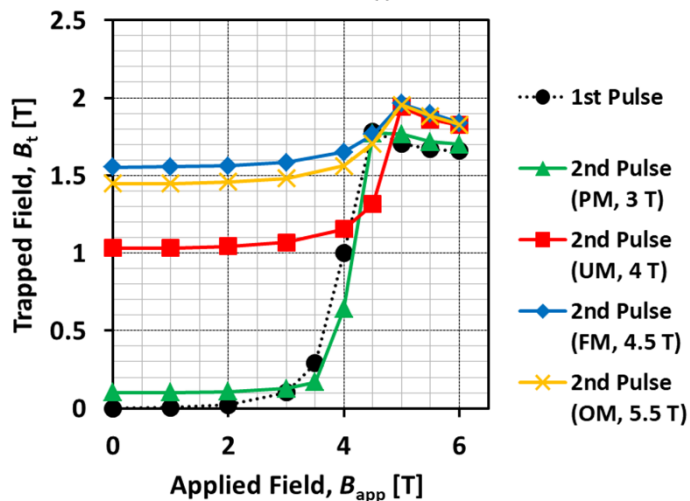
77 K



50 K



65 K



For all  $T$ , trapped field after 2<sup>nd</sup> pulse exhibits two particular characteristics:

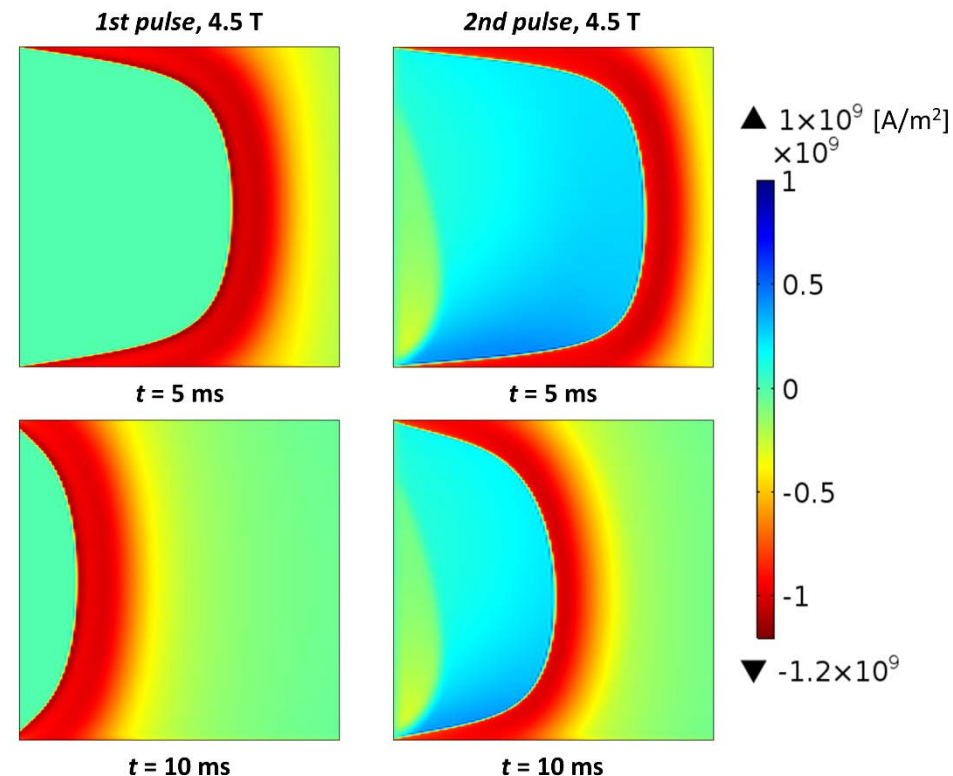
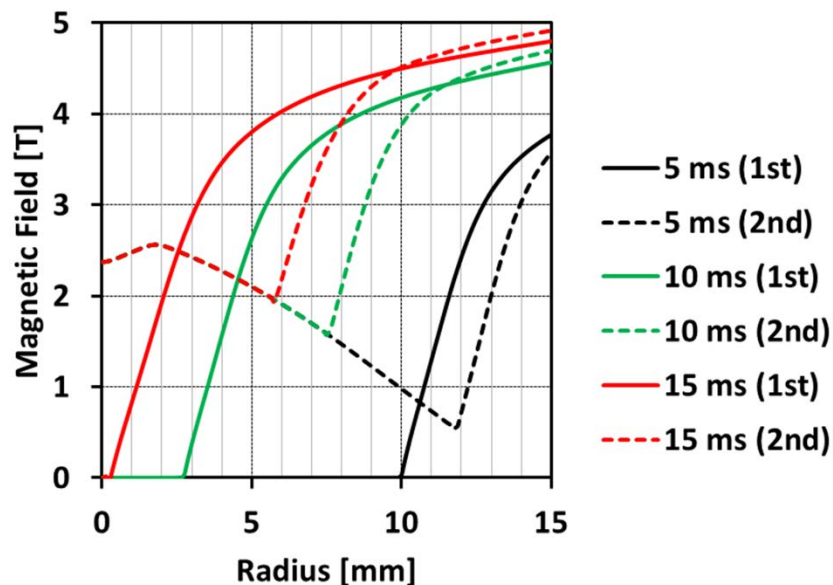
- 1) Increased trapped field,  $B_t$ , when the bulk is fully magnetised; maximum value when the 1<sup>st</sup> pulse results in full magnetisation
- 2) Increased activation field: applied field,  $B_{app}$ , required to fully magnetise the sample





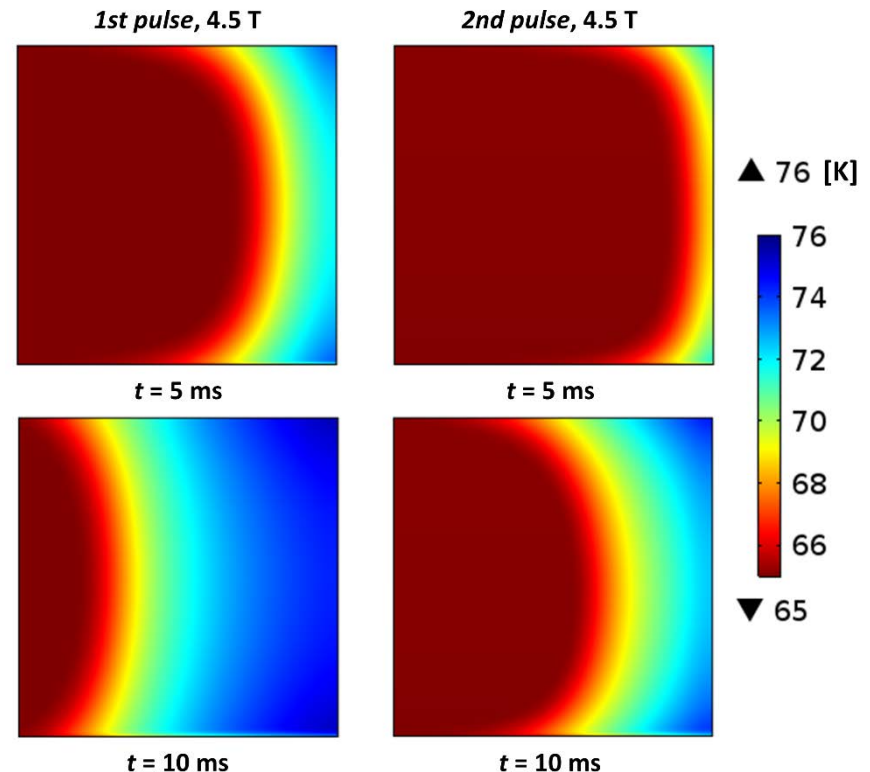
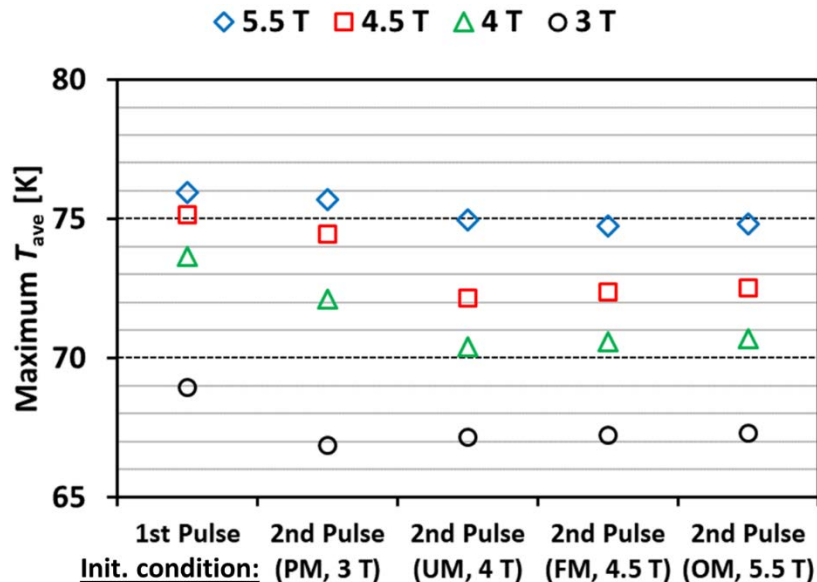
# PFM 2<sup>nd</sup> Pulse: Magnetic Flux Penetration

- Why does this occur?
  - More difficult for magnetic flux to penetrate the sample due to existing trapped field
  - Existing, induced supercurrent flows in opposite direction



# PFM 2<sup>nd</sup> Pulse: Thermal Behaviour

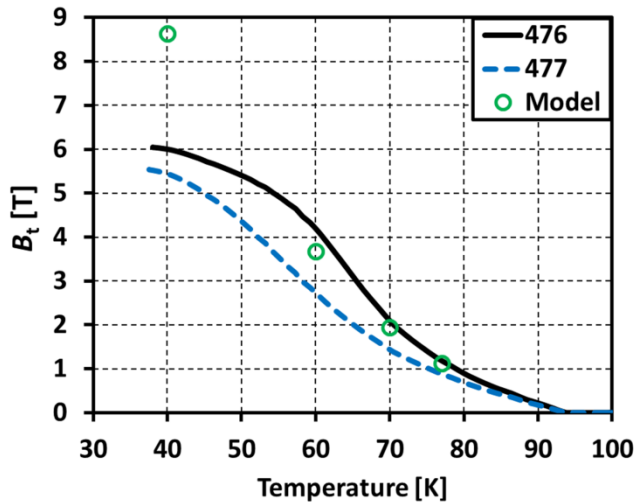
- Why does this occur?
  - Reduced dynamic movement of flux = lower temperature rise for equivalent next pulse
  - Can examine maximum average temperature,  $T_{ave,max}$ , during & after PFM:



## Case study #2:

*Split coil PFM with an iron yoke*

# Split Coil PFM with an Iron Yoke



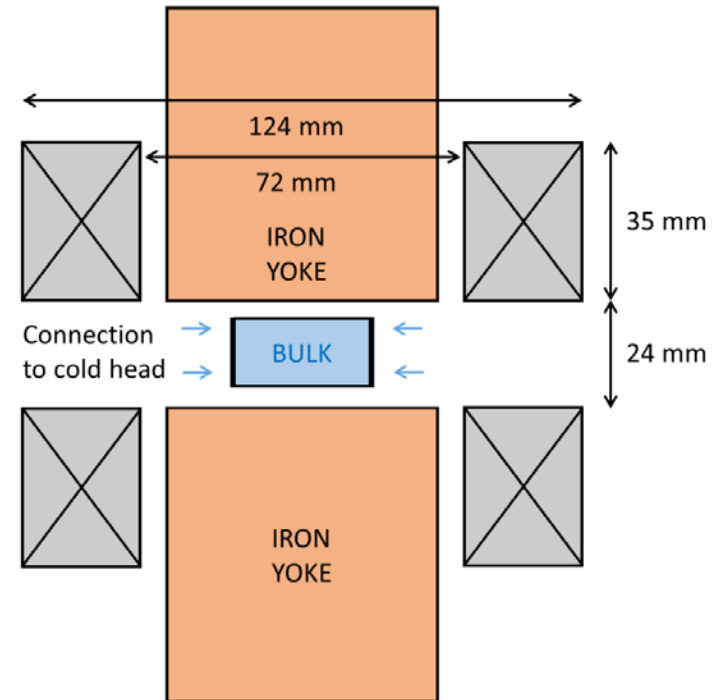
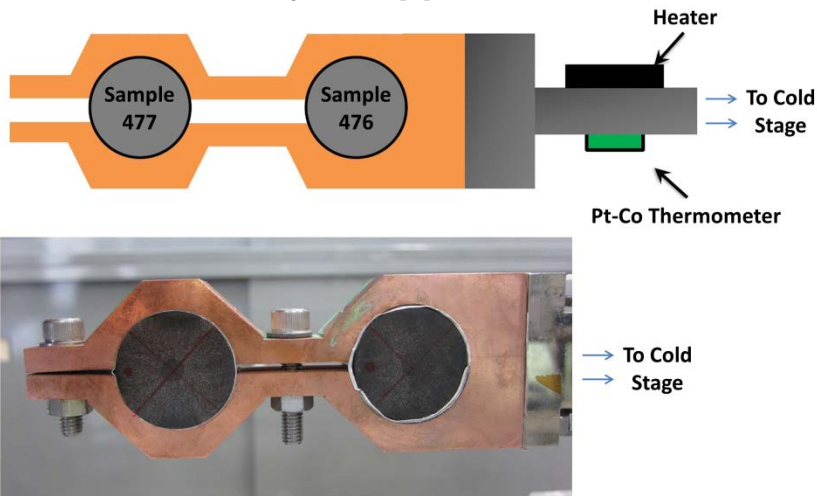
Two  $\text{GdBa}_2\text{Cu}_3\text{O}_{7-\delta}$  (15wt% Ag) samples:  
 30 mm diameter, 15 mm thickness

#476 sample  $B_t = 6 \text{ T}$  (40 K)

3.11 T (65 K)

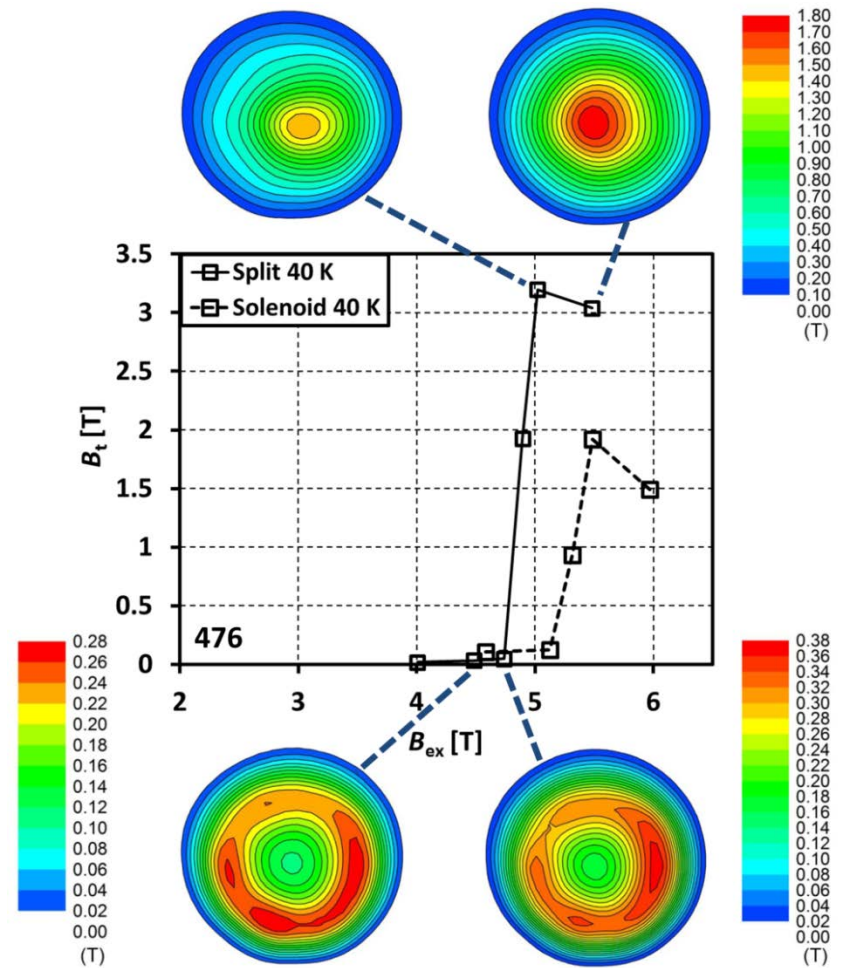
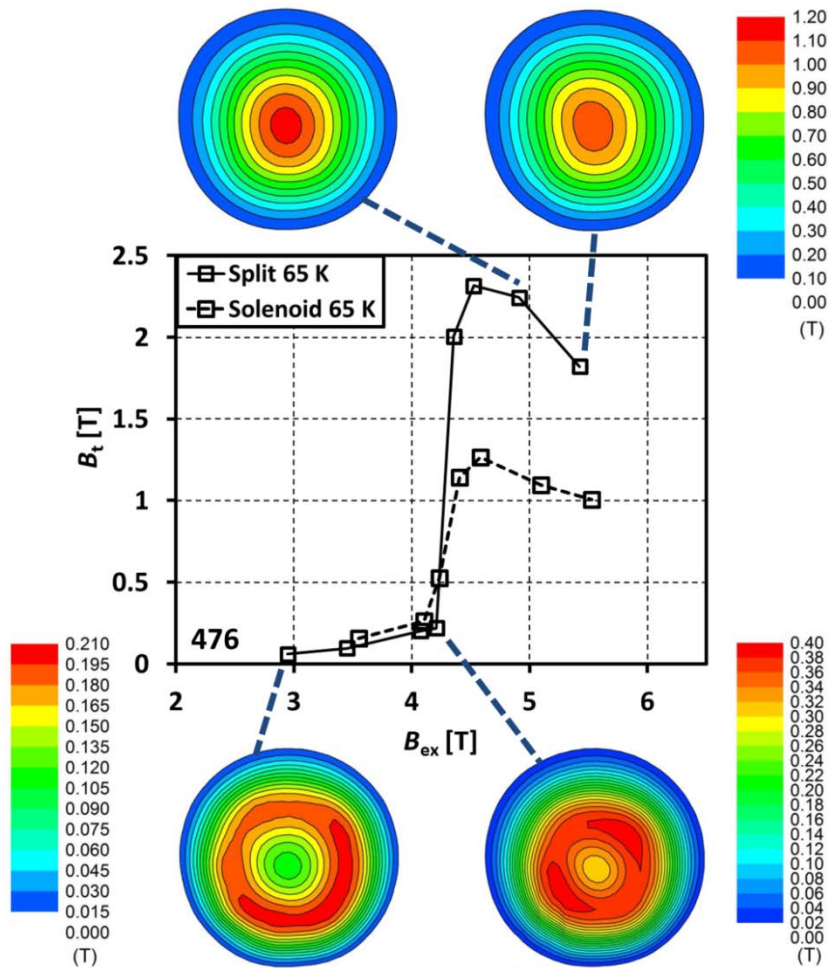
#477 sample  $B_t = 5.44 \text{ T}$  (40 K)

2.02 T (65 K)



Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **29** (2016) 074003

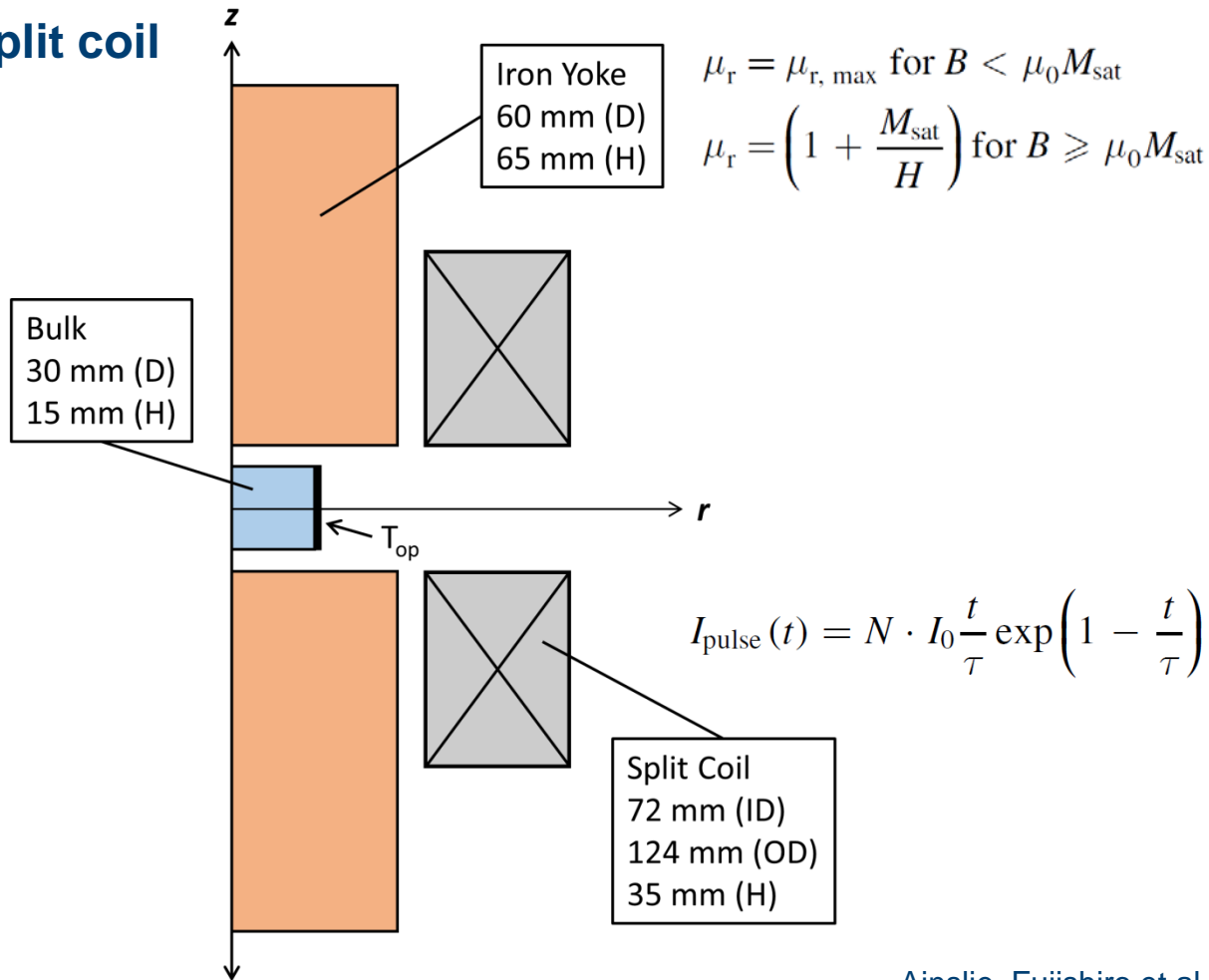
# Split Coil PFM with an Iron Yoke



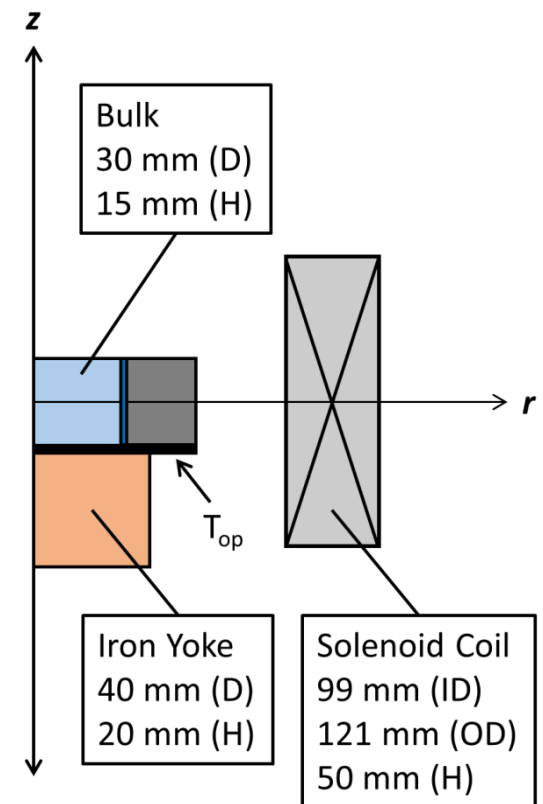
Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **29** (2016) 074003

# Split Coil PFM with an Iron Yoke

## Split coil



## Solenoid coil

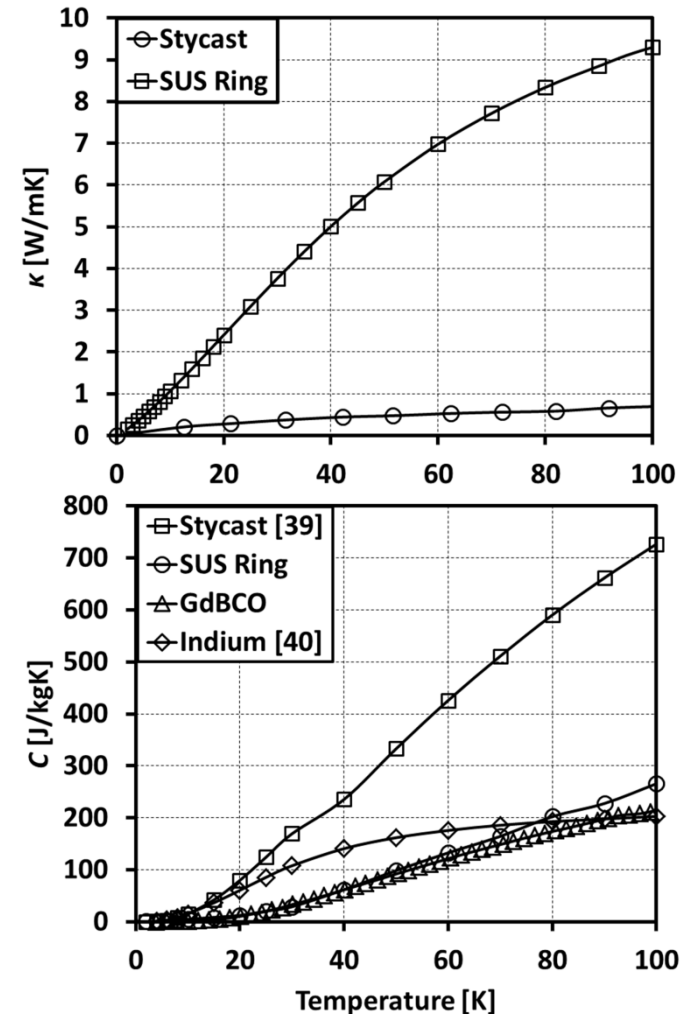
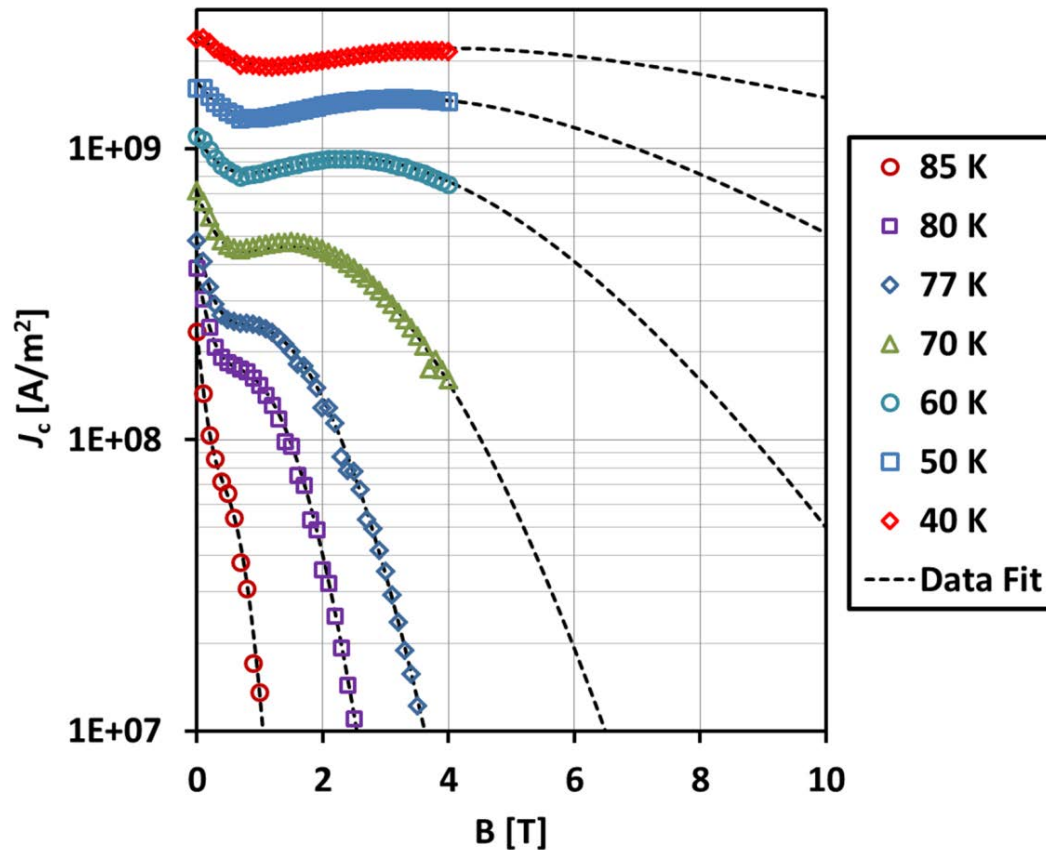


Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **29** (2016) 074003



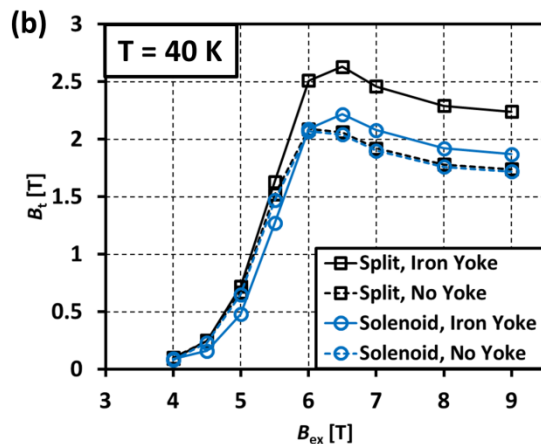
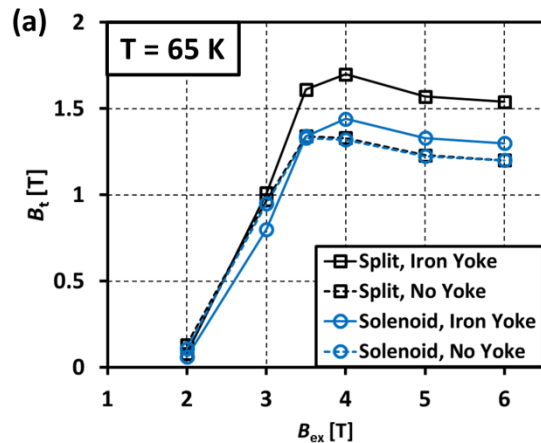
# Split Coil PFM with an Iron Yoke

$$J_c(B) = J_{c1} \exp\left(-\frac{B}{B_L}\right) + J_{c2} \frac{B}{B_{\max}} \exp\left[\frac{1}{y} \left(1 - \left(\frac{B}{B_{\max}}\right)^y\right)\right]$$

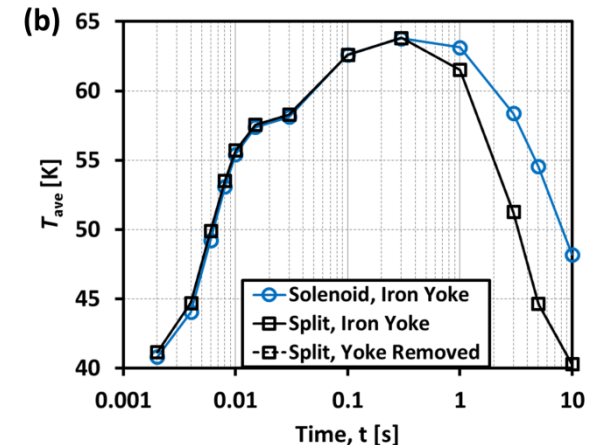
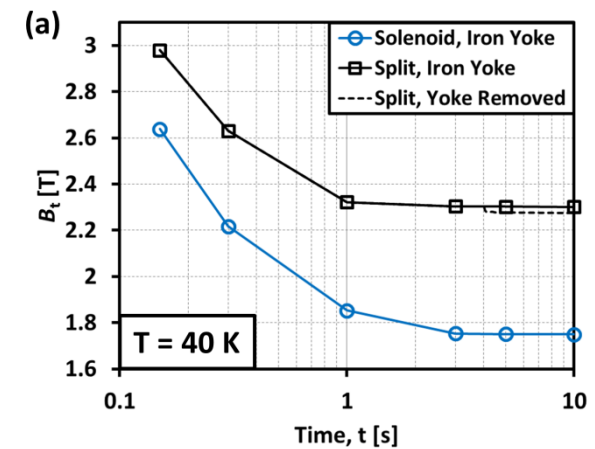
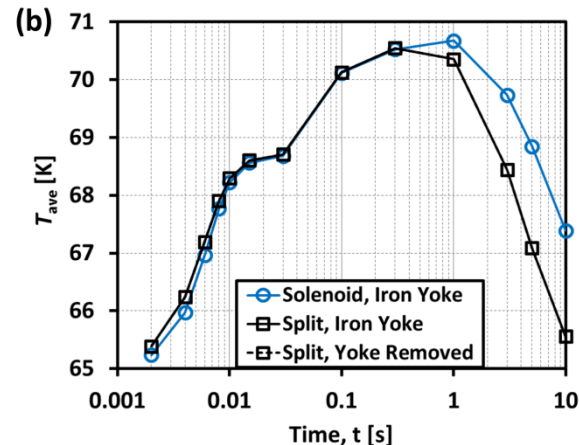
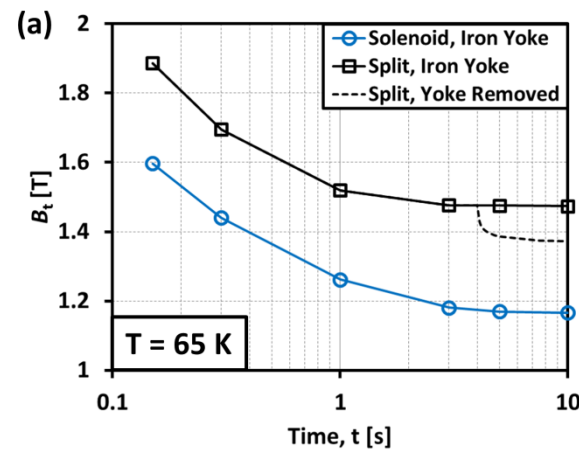


# Split Coil PFM with an Iron Yoke

## Trapped field @ $t = 300$ ms



## Trapped field & $T_{ave}$ with time, incl. yoke removed

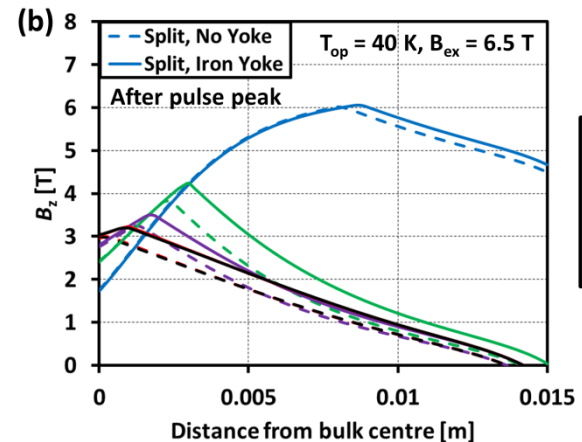
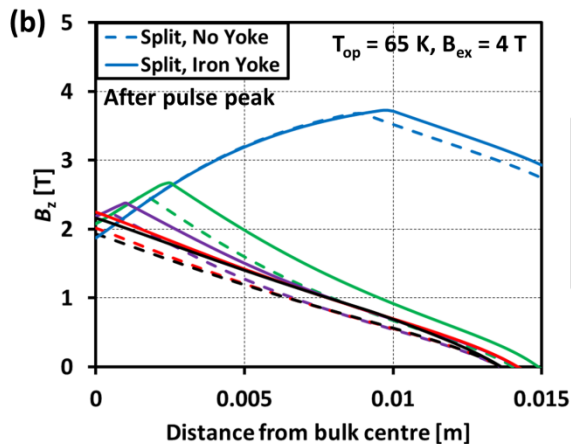
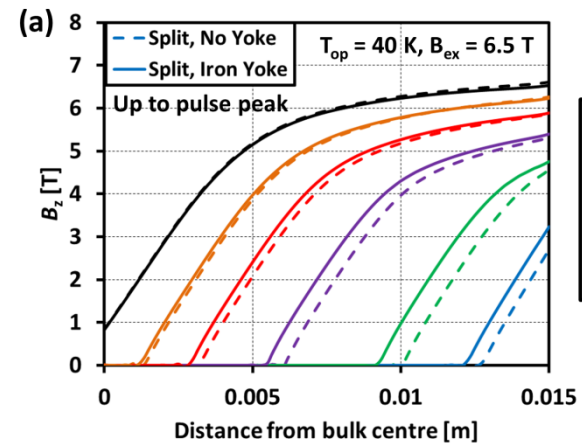
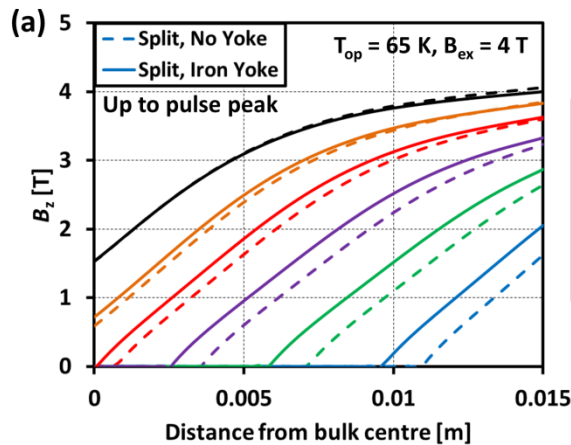
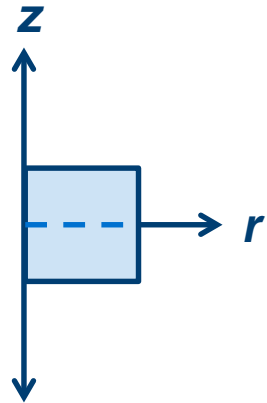


Ainslie, Fujishiro et al. *Supercond. Sci. Technol.* **29** (2016) 074003

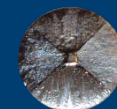


# Split Coil PFM with an Iron Yoke

## Magnetic flux entry & exit during & after pulse – split coil with & without iron yoke



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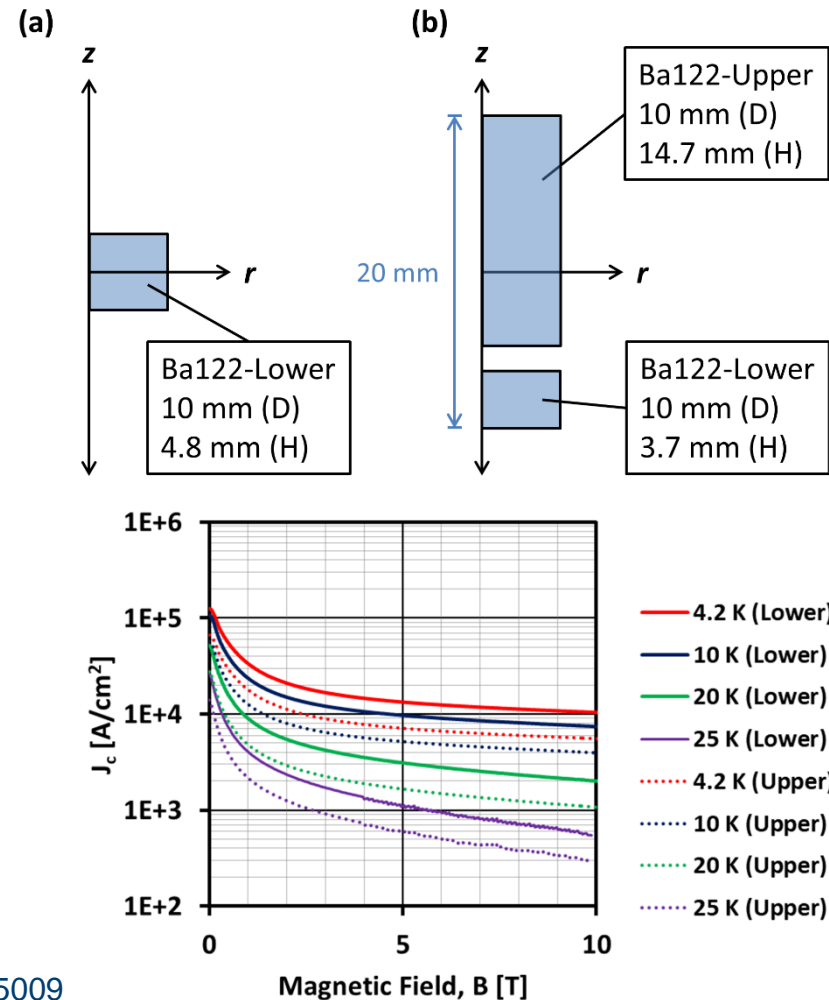


## Case study #3:

*Field cooling (FC) magnetisation of iron-pnictide (Ba122) bulks*

# Modelling Iron-Pnictide Bulks

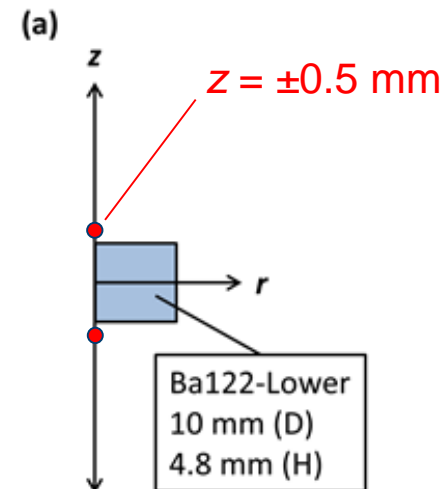
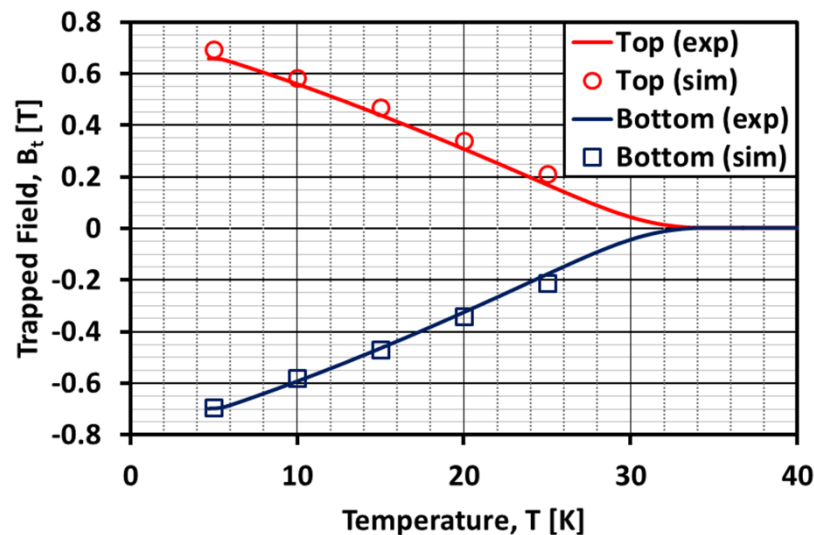
- Weiss et al. demonstrated  $> 1$  T trapped in a stack of iron-pnictide (Ba122) bulks
- Advantages:
  - Low anisotropy
  - High upper critical magnetic field ( $H_{c2}$ )
  - High homogeneity
  - Scalable, low-cost fabrication



Ainslie, Fujishiro, Yamamoto et al. *Supercond. Sci. Technol.* **30** (2017) 105009

# Modelling Iron-Pnictide Bulks

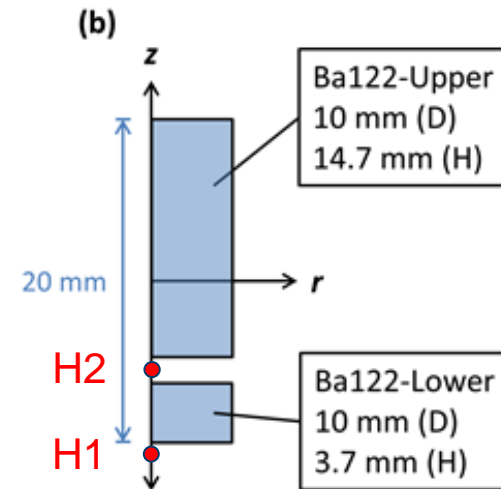
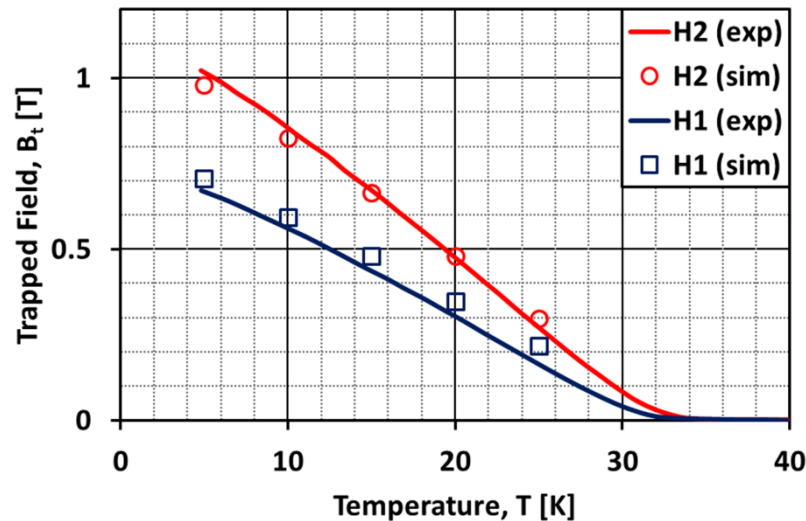
- Same ZFC model as before  
→ approximates FC, applied field  $\gg$  full penetration field
- $n = 50$   
→ observed flux creep much weaker than HTS bulks



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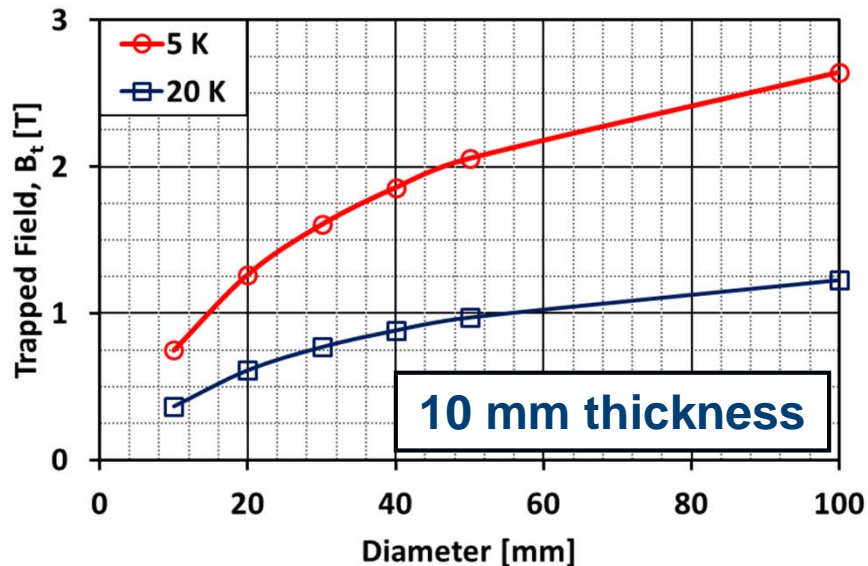


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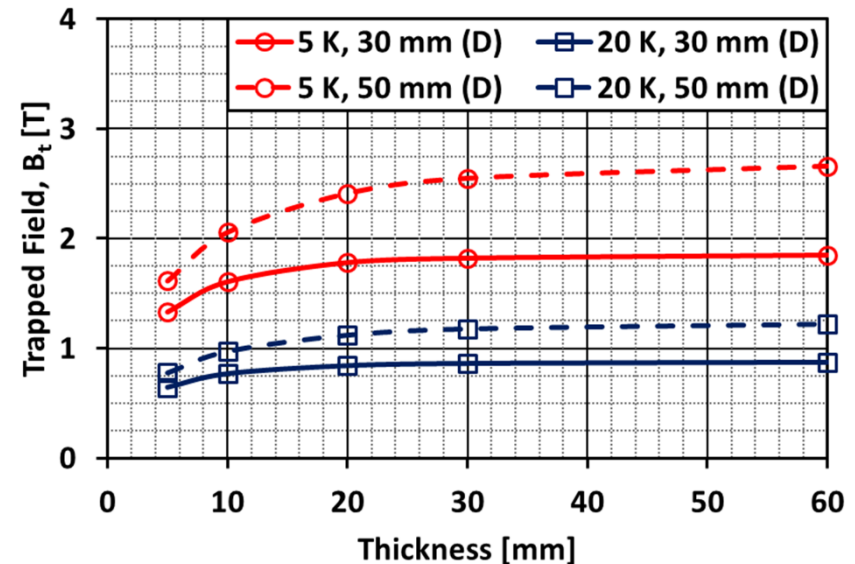
# Modelling Iron-Pnictide Bulks

Influence of geometric parameters can be predicted:

Diameter dependence



Thickness dependence



- With current state-of-the-art properties,  $> 2$  T @ 5 K ( $> 1$  T @ 20 K) for  $D = 50$  mm
- Appropriate aspect ratio 1-1.5 (radius : thickness)

Ainslie, Fujishiro, Yamamoto et al. *Supercond. Sci. Technol.* **30** (2017) 105009

# Summary

- **Overview of numerical modelling of magnetisation**
- **Case studies:**
  - Multi-pulse, pulsed field magnetisation (PFM) of bulk high-temperature superconductors
  - Split coil PFM with an iron yoke
  - Field cooling (FC) magnetisation of iron-pnictide (Ba122) bulks

# Thank you for listening