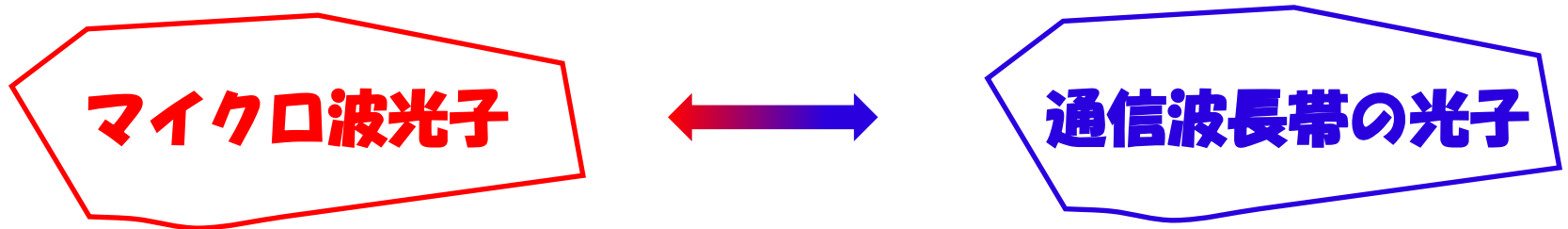


量子トランスデューサー

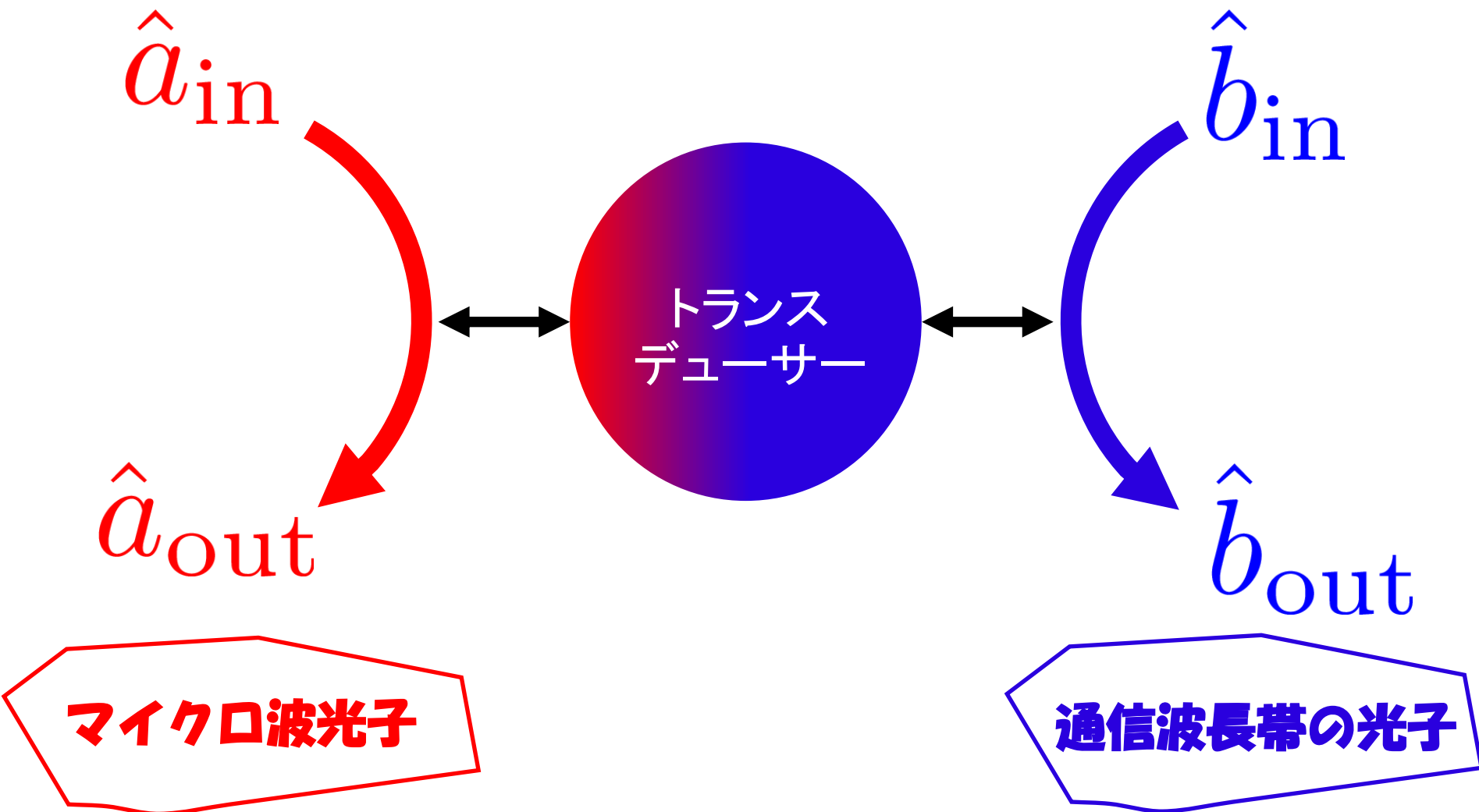
Quantum Transducers



東京大学先端科学技術研究センター
宇佐見康二

RCAST, UTokyo
Koji Usami

量子トランスデューサーとは？



交換関係

$\hat{a}_{\text{in}}$

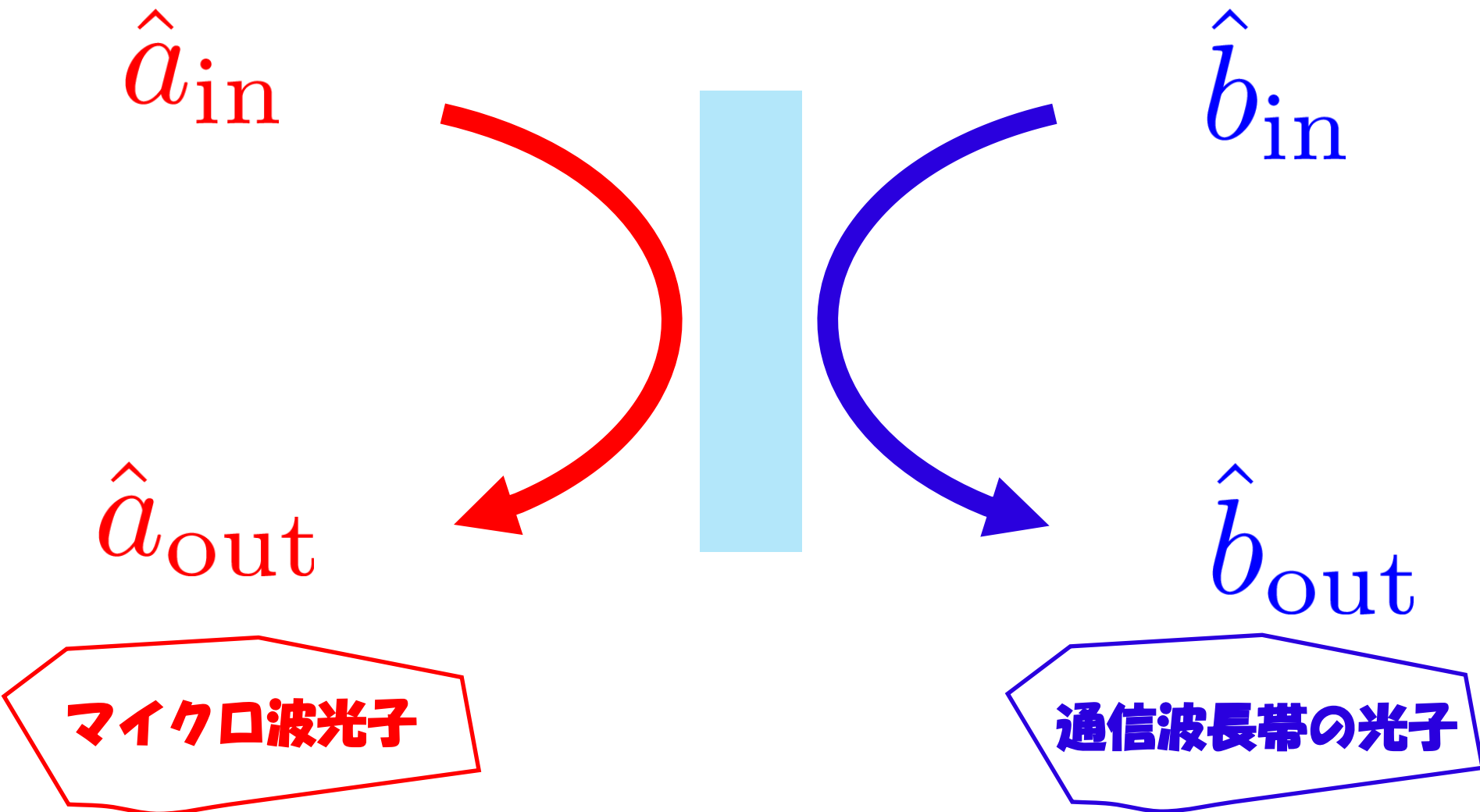
$\hat{b}_{\text{in}}$

$\hat{a}_{\text{out}}$

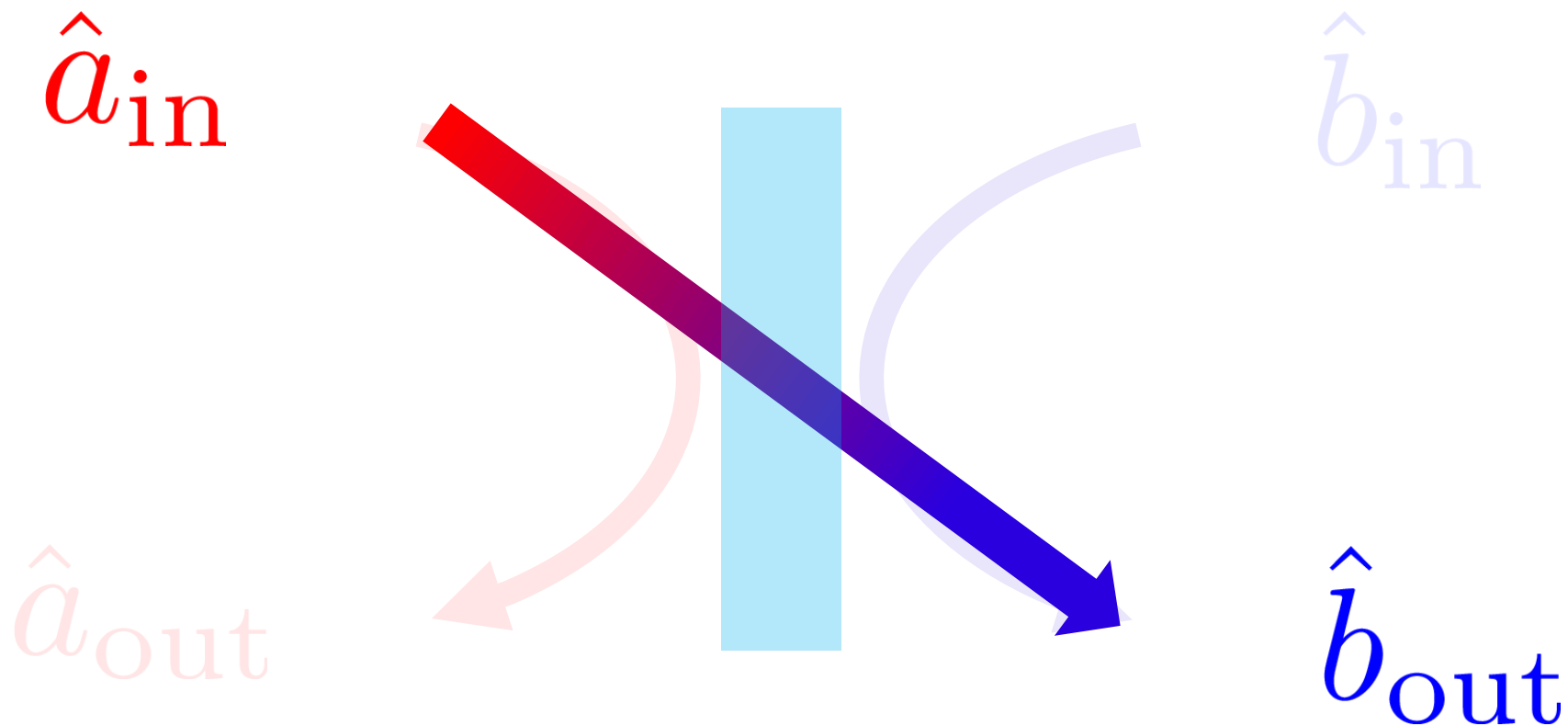
$\hat{b}_{\text{out}}$

マイクロ波光子

通信波長帯の光子



交換關係



$$\left[\hat{a}_{\text{in}}(t), \hat{a}_{\text{in}}^{\dagger}(t') \right] = \delta(t - t')$$

$$\left[\hat{b}_{\text{out}}(t), \hat{b}_{\text{out}}^{\dagger}(t') \right] = \delta(t - t')$$

交換關係

損失: $\hat{b}_{\text{out}} = \eta \hat{a}_{\text{in}}$ ($\eta \leq 1$)

增幅: $\hat{b}_{\text{out}} = G \hat{a}_{\text{in}}$ ($G \geq 1$)

$$[\hat{a}_{\text{in}}(t), \hat{a}_{\text{in}}^\dagger(t')] = \delta(t - t')$$

$$[\hat{b}_{\text{out}}(t), \hat{b}_{\text{out}}^\dagger(t')] \neq \delta(t - t')$$

交換関係

損失: $\hat{b}_{\text{out}} = \eta \hat{a}_{\text{in}} + \sqrt{1 - \eta^2} \hat{c}_{\text{in}} \quad (\eta \leq 1)$

増幅: $\hat{b}_{\text{out}} = G \hat{a}_{\text{in}} + \sqrt{G^2 - 1} \hat{d}_{\text{in}} \quad (G \geq 1)$

ノイズ項

$$[\hat{a}_{\text{in}}(t), \hat{a}_{\text{in}}^\dagger(t')] = \delta(t - t')$$

$$[\hat{b}_{\text{out}}(t), \hat{b}_{\text{out}}^\dagger(t')] = \delta(t - t')$$

交換関係

損失: $\hat{b}_{\text{out}} = \eta \hat{a}_{\text{in}}$ 損失なし ($\eta = 1$)

増幅: $\hat{b}_{\text{out}} = G \hat{a}_{\text{in}}$ 増幅なし ($G = 1$)

$$[\hat{a}_{\text{in}}(t), \hat{a}_{\text{in}}^\dagger(t')] = \delta(t - t')$$

$$[\hat{b}_{\text{out}}(t), \hat{b}_{\text{out}}^\dagger(t')] = \delta(t - t')$$

ユニティゲインパラメトリック増幅器！

$$\Omega_p = \Omega - \omega$$

ポンプ光子

$$\hat{b}_{\text{out}} \longleftrightarrow \hat{a}_{\text{in}}$$

通信波長帯の光子

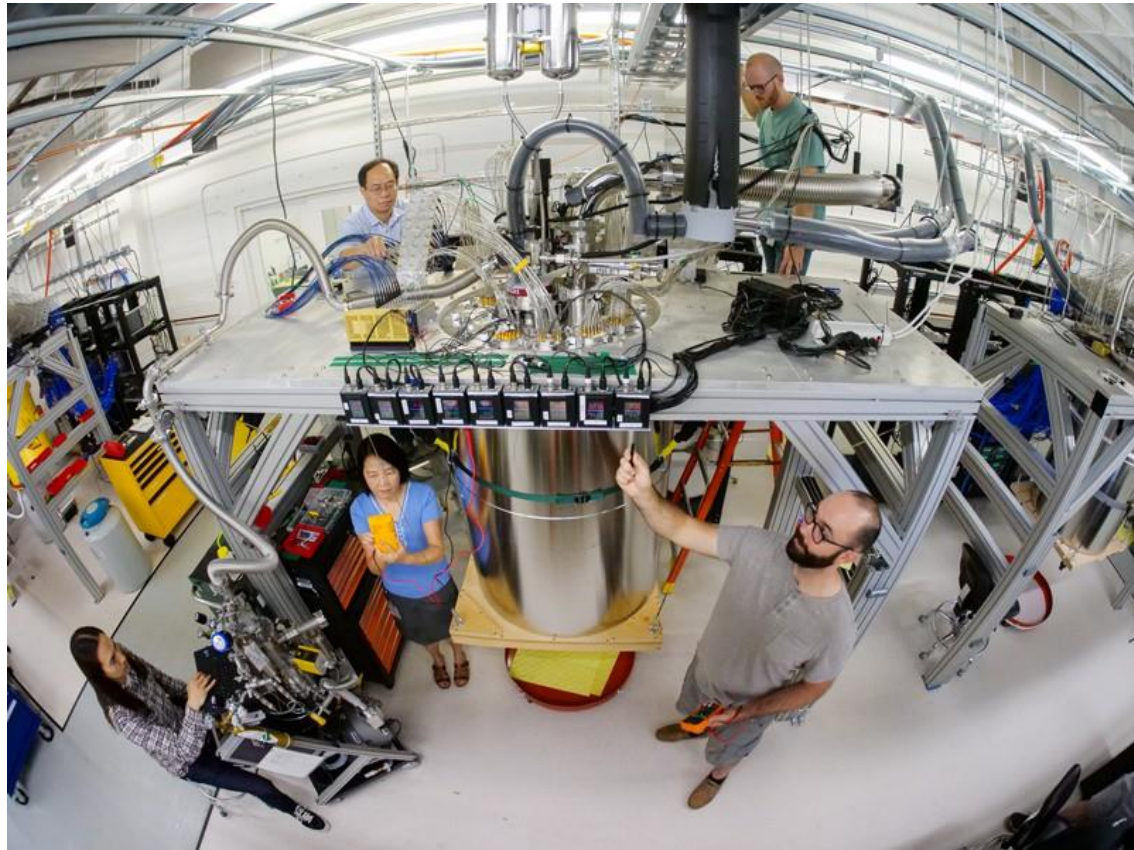
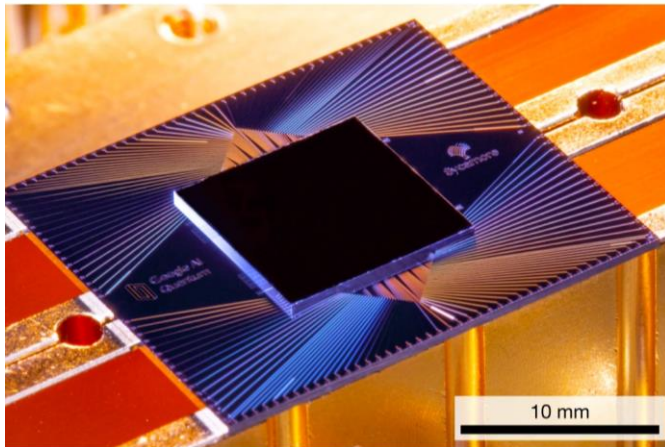
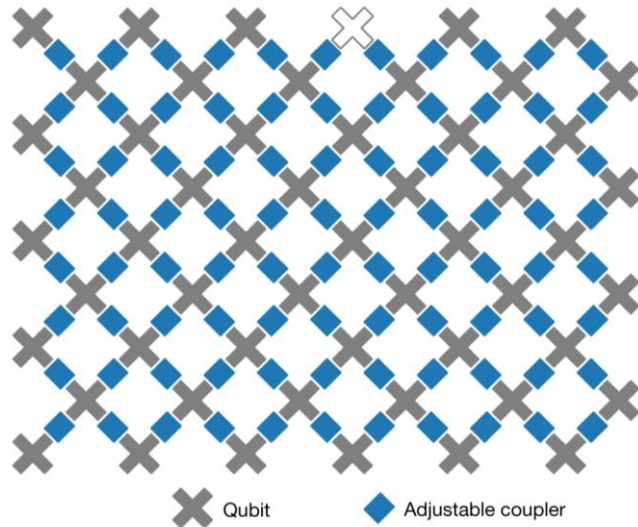
マイクロ波光子

$$\Omega/2\pi \sim 200 \text{ THz} - 600 \text{ THz} \quad \omega/2\pi \sim 0.1 \text{ MHz} - 10 \text{ GHz}$$

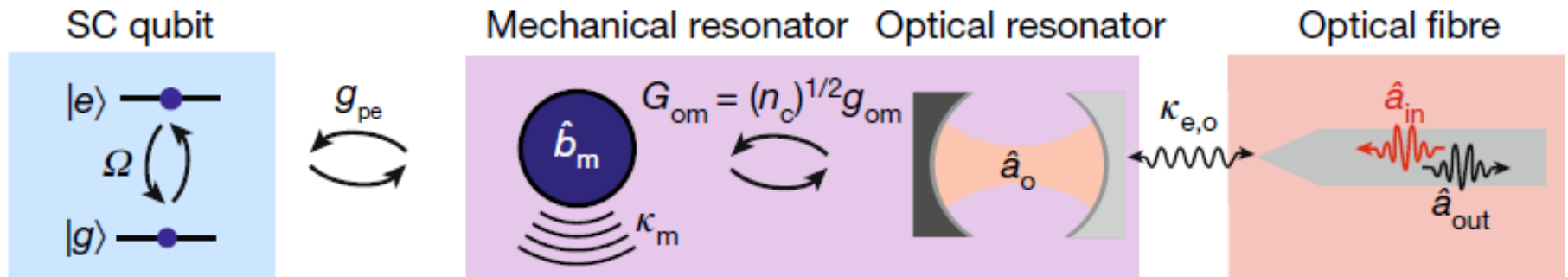
動機

- 量子ビットのネットワーク化
- 熱負荷の緩和
- 微弱ラジオ波信号の光検出

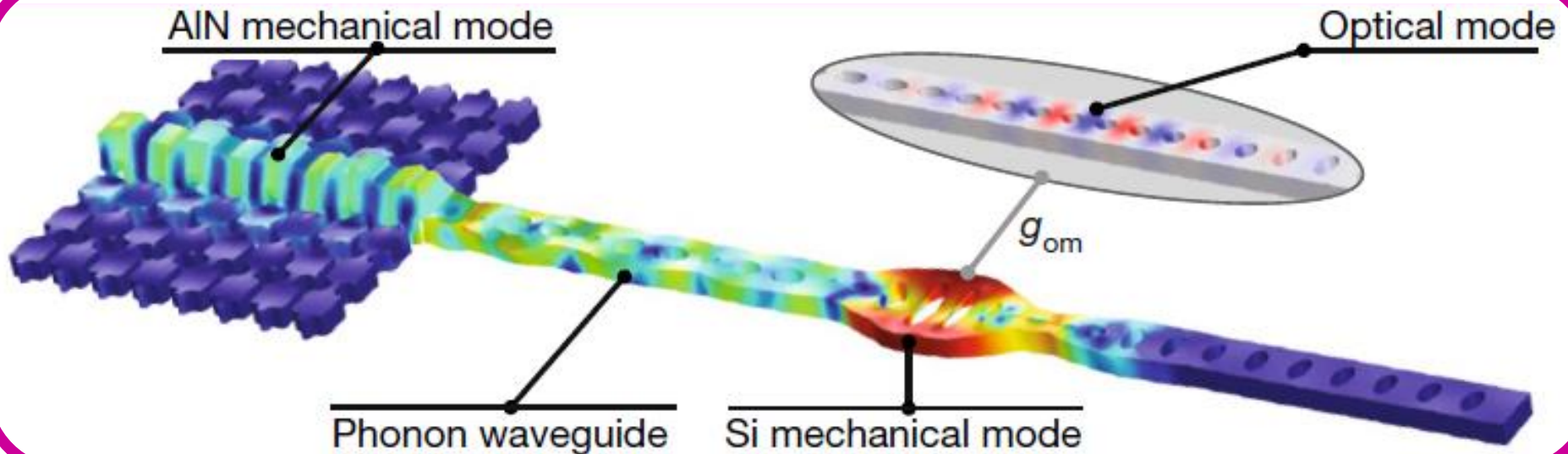
動機 1：量子ビットのネットワーク化



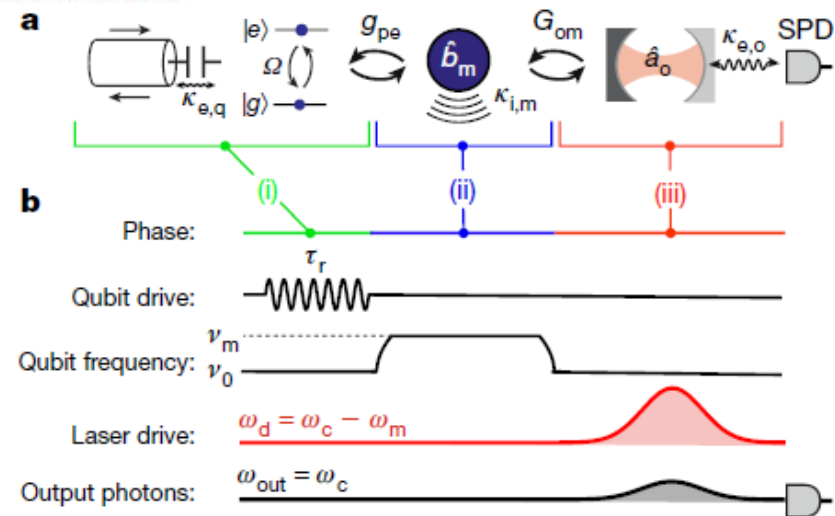
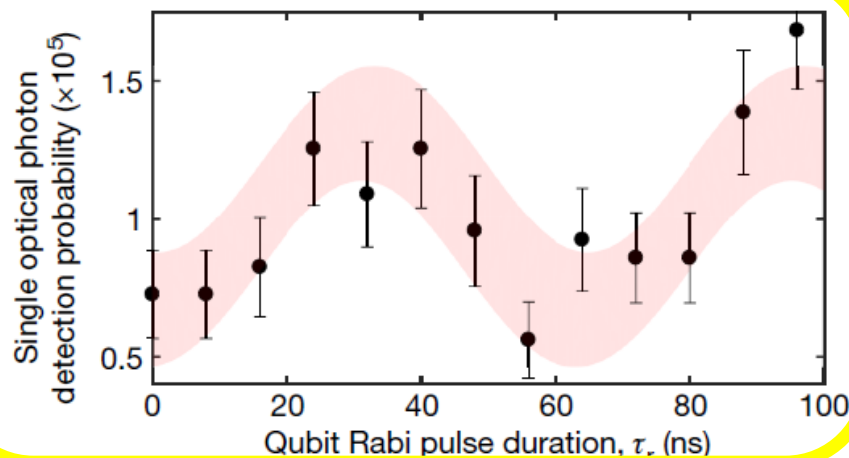
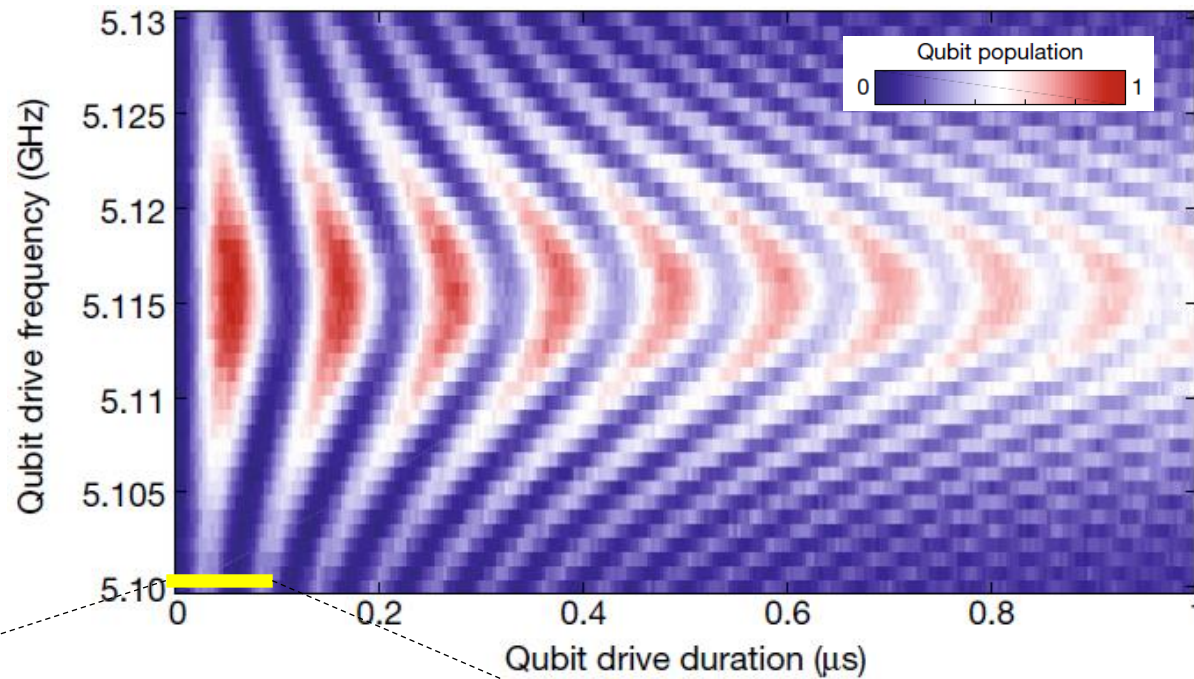
ラビ振動の光学的検出



Mechanical “super-mode”



ラビ振動の光学的検出

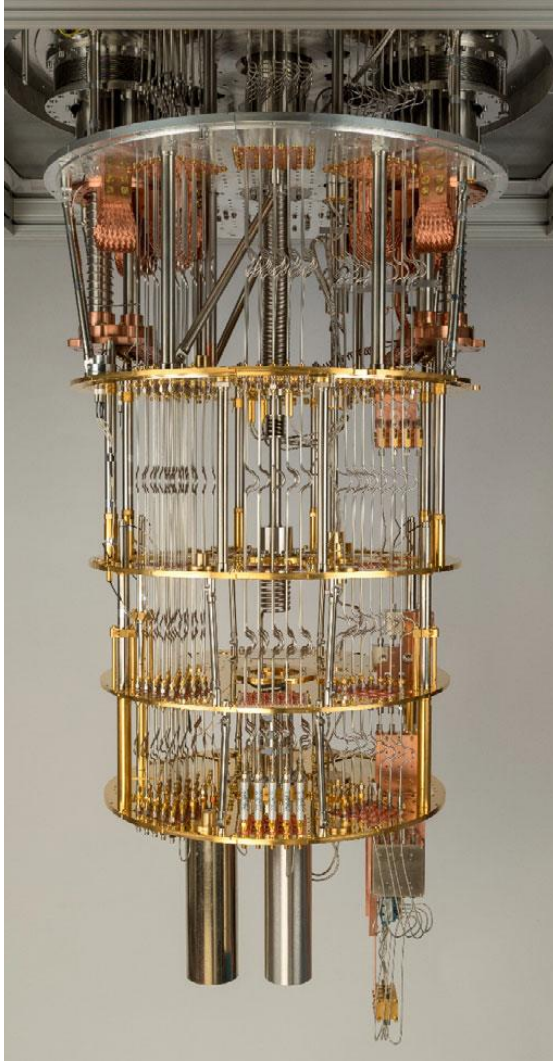


M. Mirhosseini *et al.*, Nature **588**, 599 (2020)

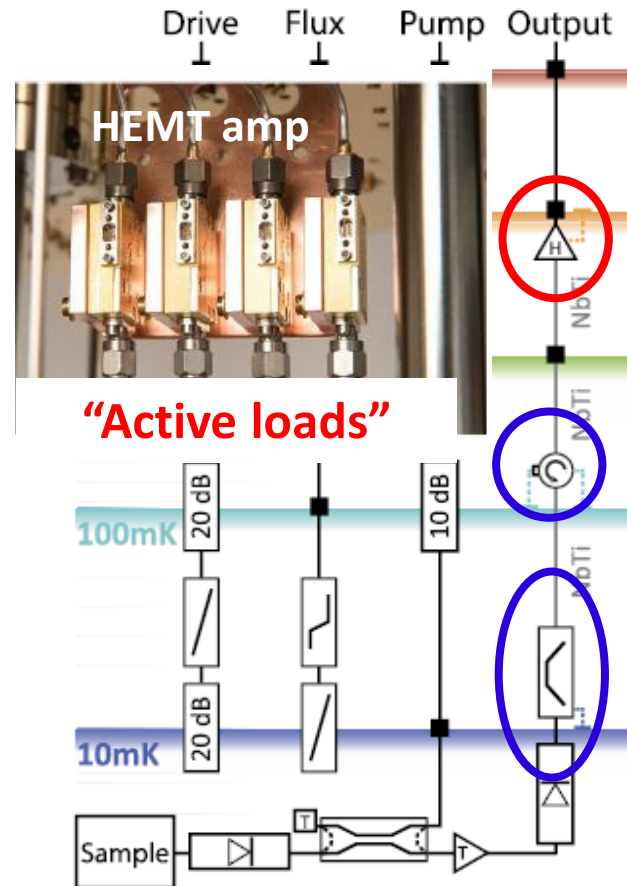
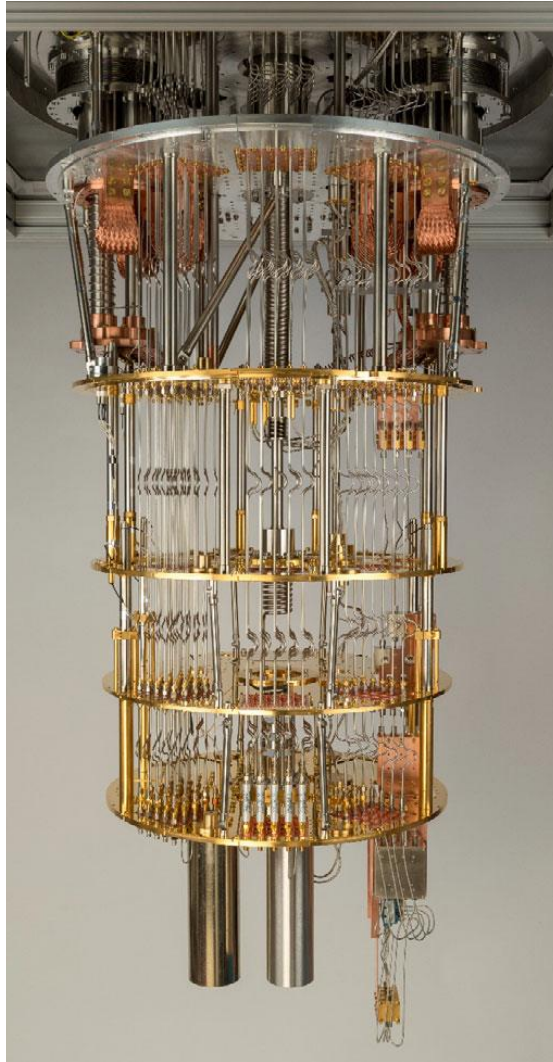
動機

- 量子ビットのネットワーク化
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- 微弱ラジオ波信号の光検出

動機 2 : 熱負荷の緩和



動機 2 : 熱負荷の緩和

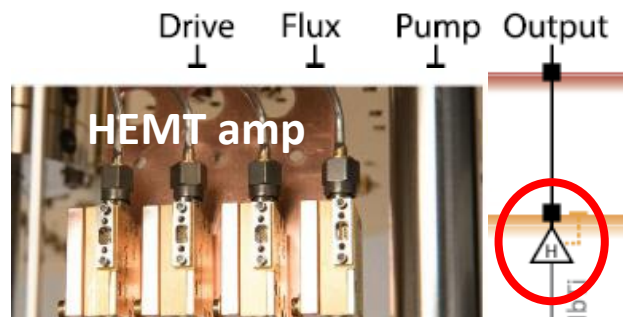
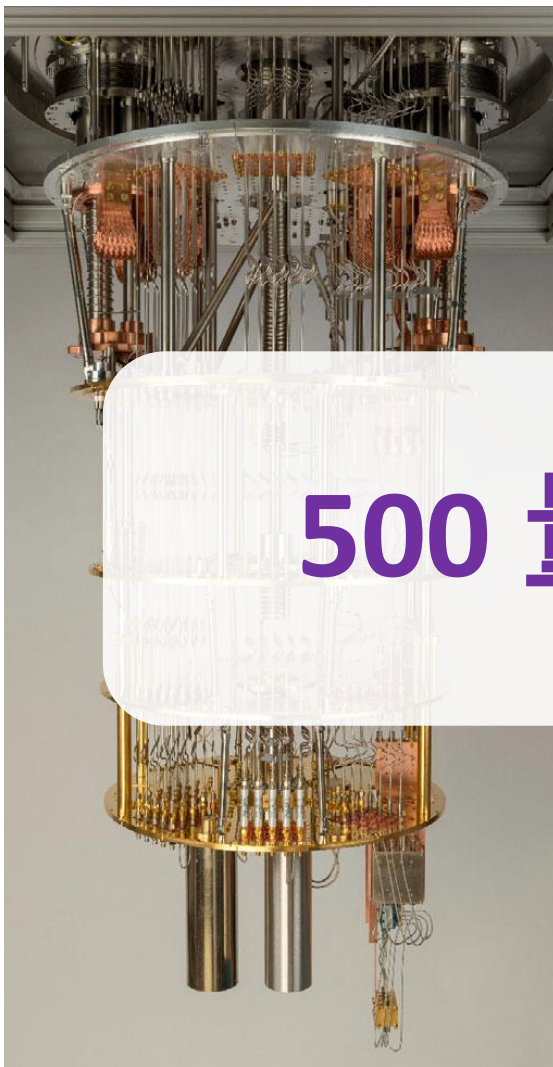


“Passive loads”

Circulators

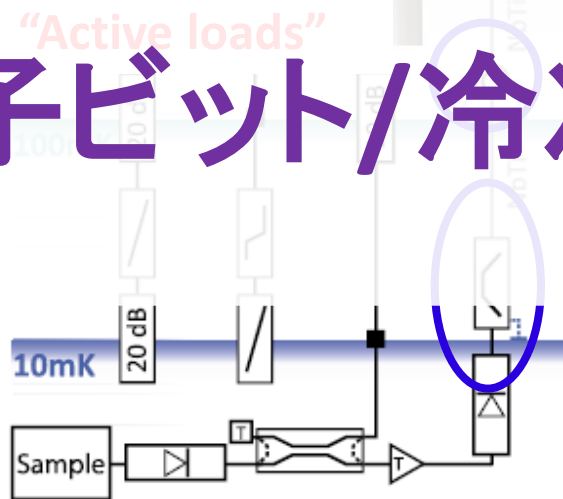
Band-pass filters

動機 2：熱負荷の緩和

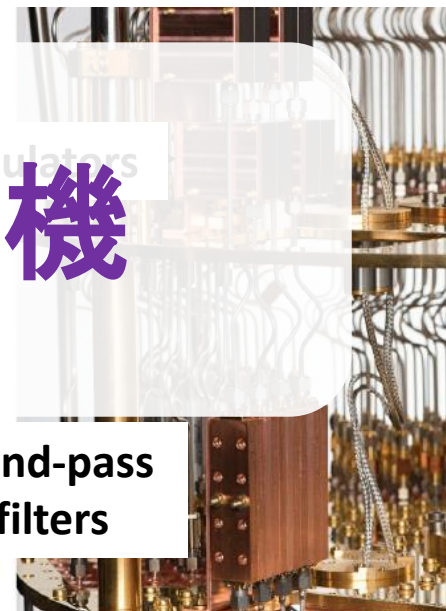


“Passive loads”

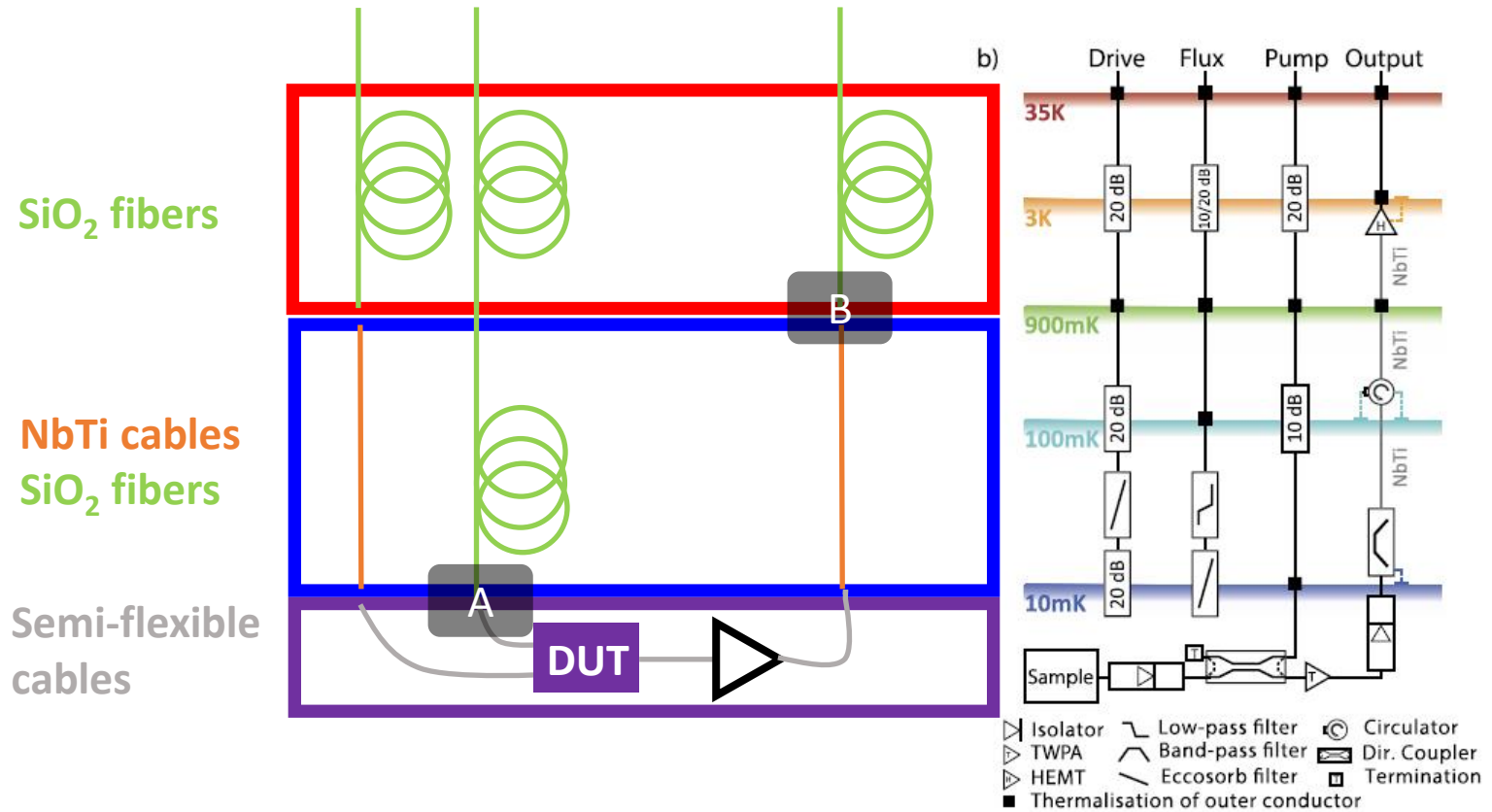
500 量子ビット/冷凍機



Band-pass filters



イントラフリッジ光リンク



Control and readout of a superconducting qubit using a photonic link

<https://doi.org/10.1038/s41586-021-03268-x>

Received: 3 August 2020

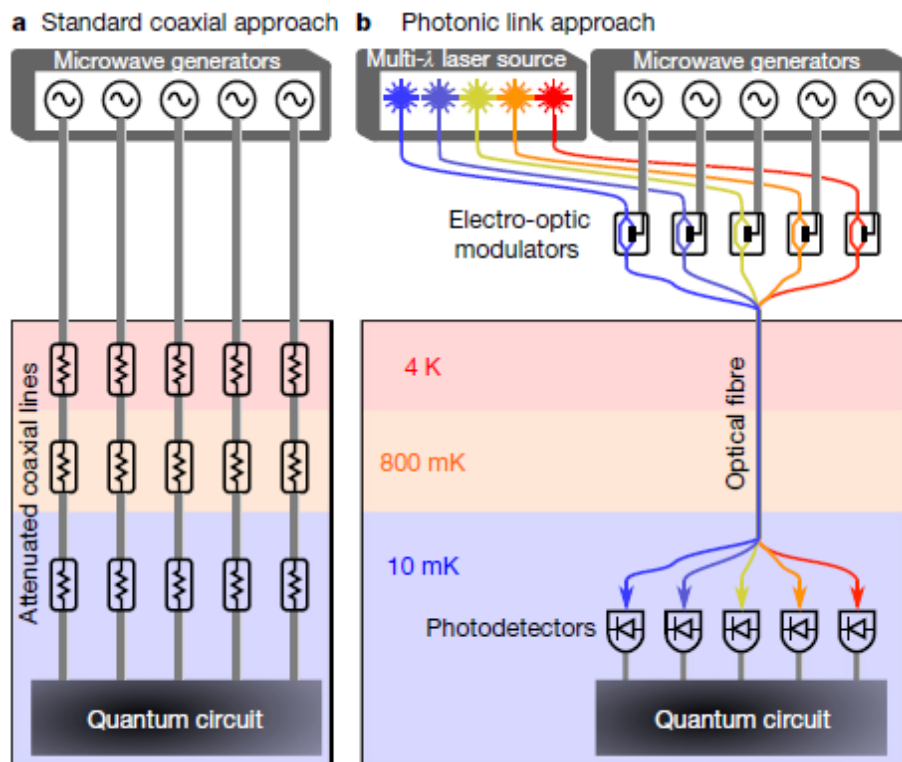
Accepted: 20 January 2021

Published online: 24 March 2021

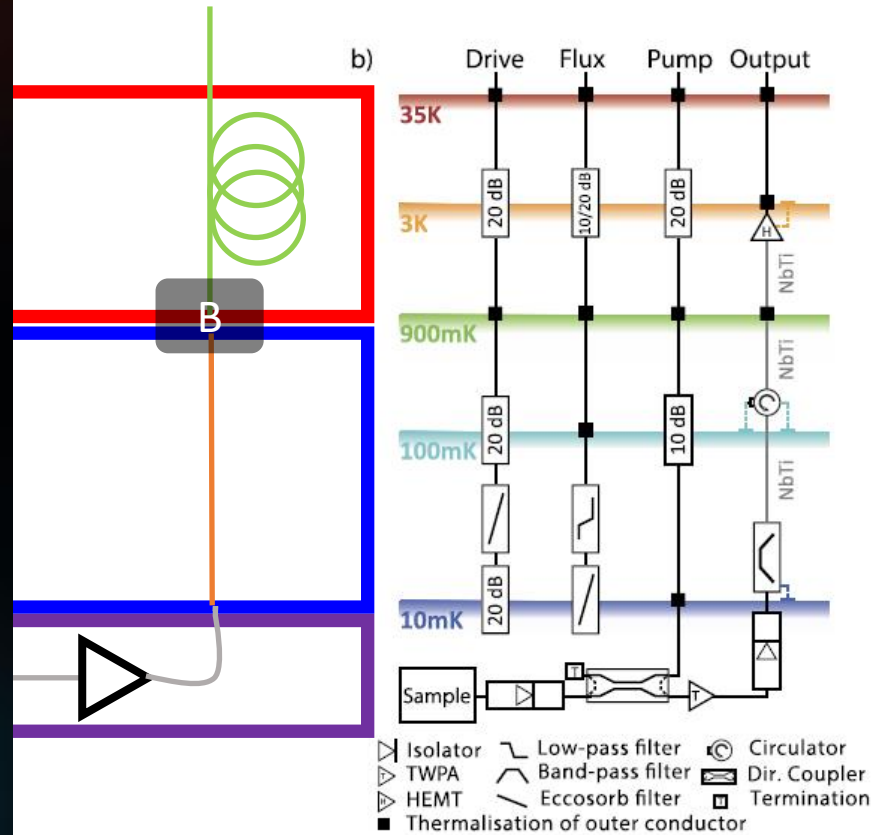
 Check for updates

F. Lecocq^{1,2}, F. Quinlan¹, K. Cicak¹, J. Aumentado¹, S. A. Diddams^{1,2} & J. D. Teufel¹

Delivering on the revolutionary promise of a universal quantum computer will require processors with millions of quantum bits (qubits)^{1–3}. In superconducting quantum processors⁴, each qubit is individually addressed with microwave signal lines that connect room-temperature electronics to the cryogenic environment of the quantum circuit. The complexity and heat load associated with the multiple coaxial lines per qubit limits the maximum possible size of a processor to a few thousand qubits⁵. Here we introduce a photonic link using an optical fibre to guide modulated laser light from room temperature to a cryogenic photodetector⁶, capable of delivering shot-noise-limited microwave signals directly at millikelvin temperatures. By demonstrating high-fidelity control and readout of a superconducting qubit, we show that this photonic link can meet the stringent requirements of superconducting quantum information processing⁷. Leveraging the low thermal conductivity and large intrinsic bandwidth of optical fibre enables the efficient and massively multiplexed delivery of coherent microwave control pulses, providing a path towards a million-qubit universal quantum computer.



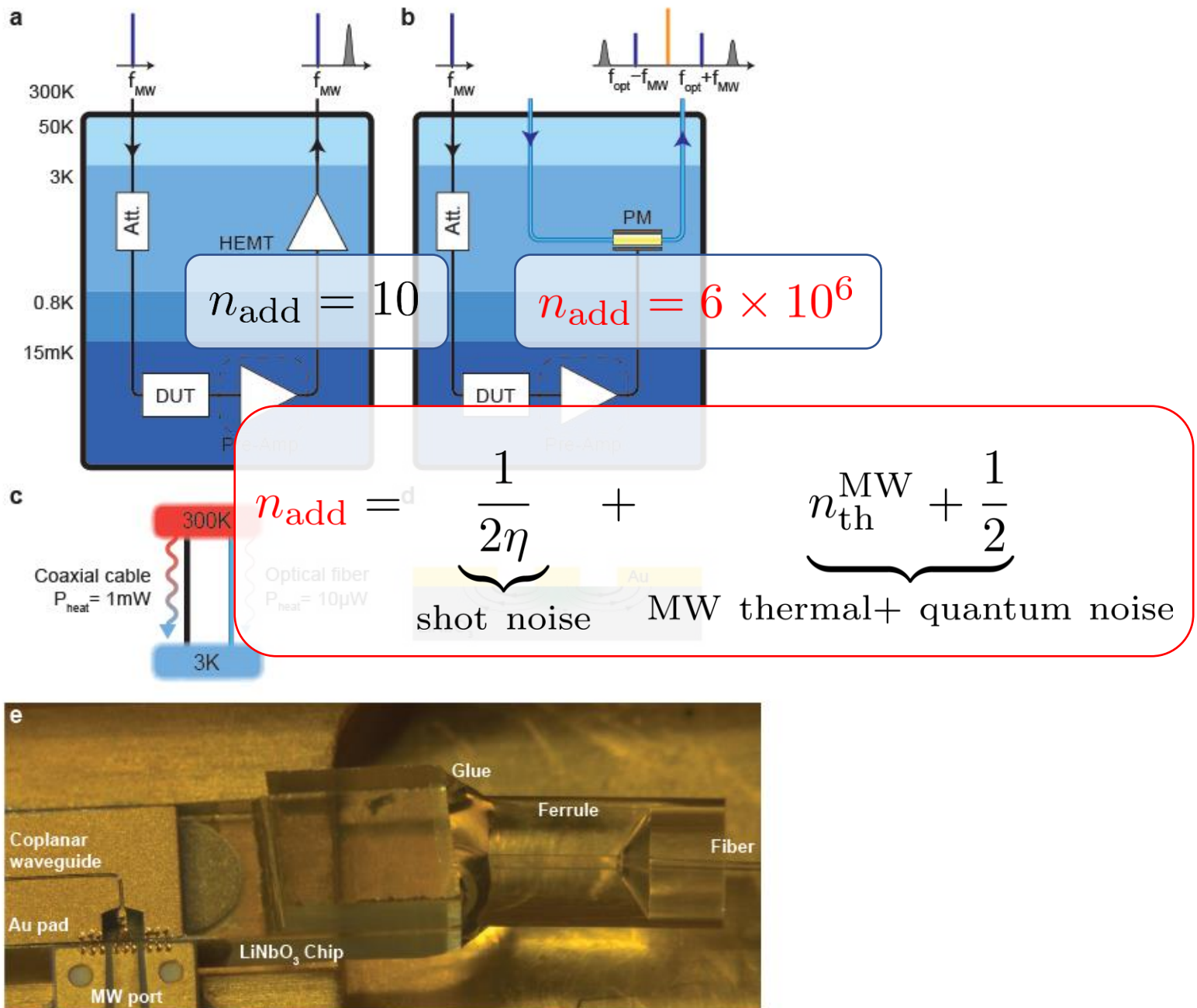
イントラフリッジ光リンク



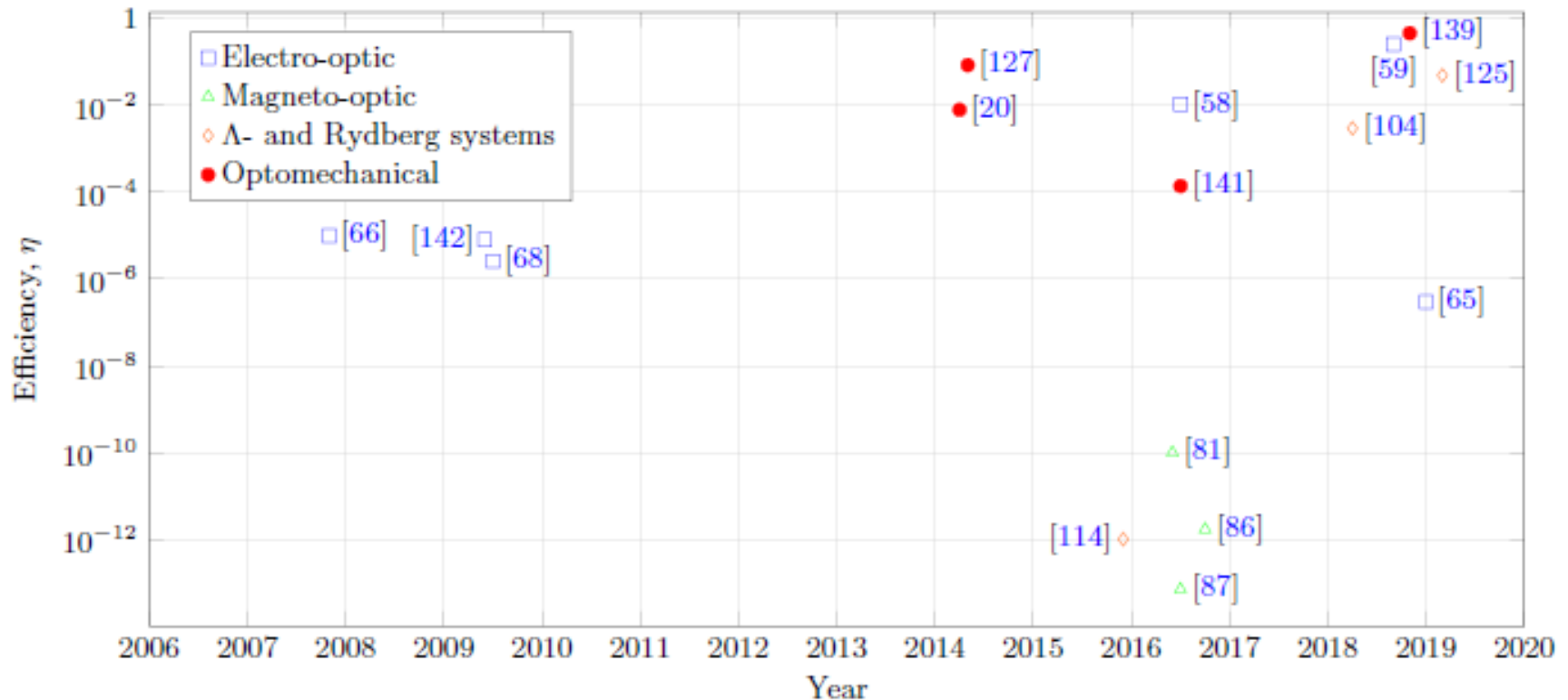
A. Youssefi et al., Nature Electronics **4**, 326 (2021)

(News & Views: K. Usami and Y. Nakamura)

B



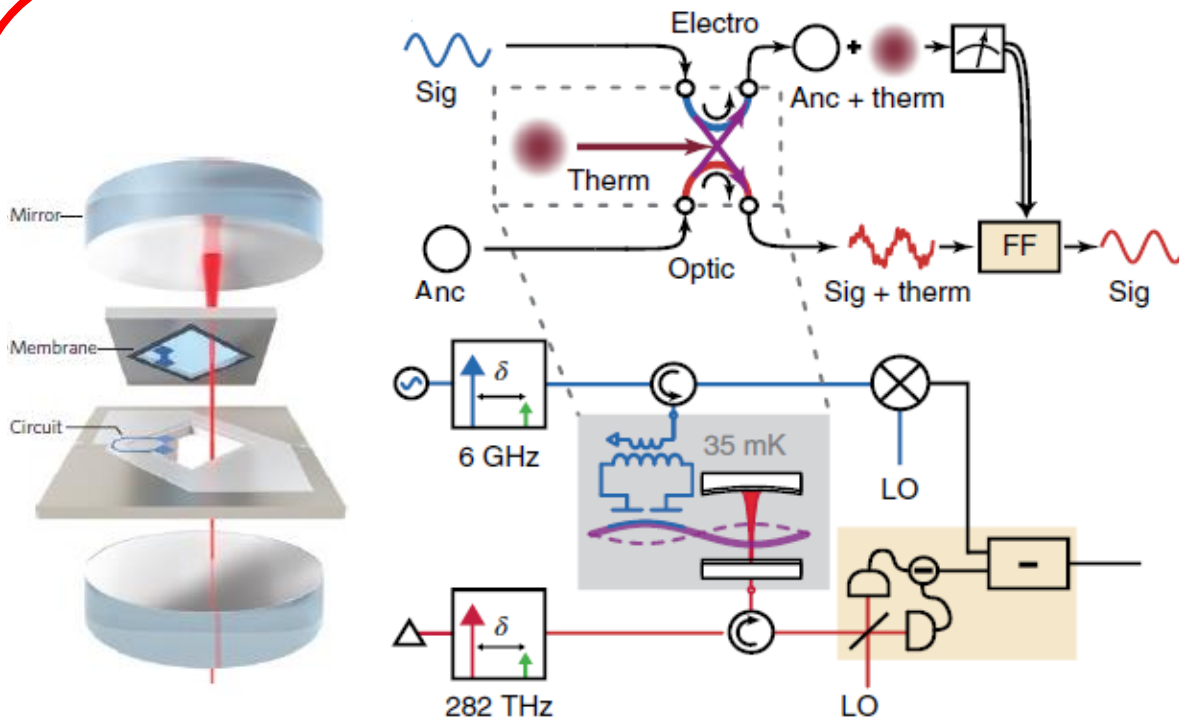
変換効率 η の変遷



Ref.	Type	Material system	Ω (GHz)	η	Bandwidth (MHz)
[58]	Electro-optic	LiNbO ₃	8.9	0.0109	1.38
[143]	Λ -system	ER:YSO	5.18	1.2×10^{-5}	0.13
[59]	Electro-optic	AlN	8.31	0.259	0.59
[81]	Magneto-optic	YIG	10.5	10^{-10}	2.28
[125]	Λ -system	Rubidium atoms	84	5×10^{-2}	15
[139]	Optomechanical	Si ₃ N ₄	6.16	0.47	1.2×10^{-2}

N. J. Lambert *et al.*, Adv. Quantum Technol. 3, 1900077 (2020).

変換効率 η の変遷



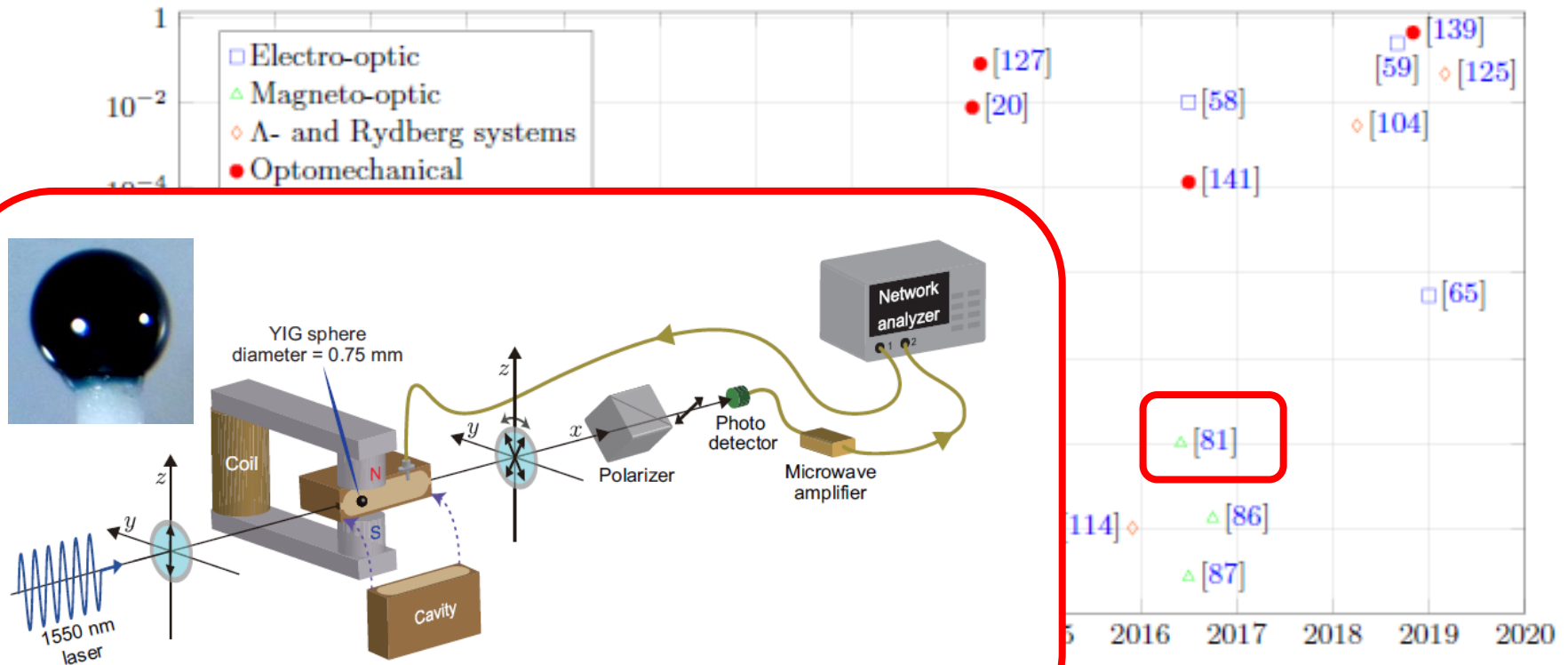
A. P. Higginbotham *et al.*, Nat. Phys. **14**, 1038 (2018)



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N. J. Lambert *et al.*, Adv. Quantum Technol. **3**, 1900077 (2020).

変換効率 η の変遷



R. Hisatomi *et al.*, Phys. Rev. B **93**, 174427 (2016)

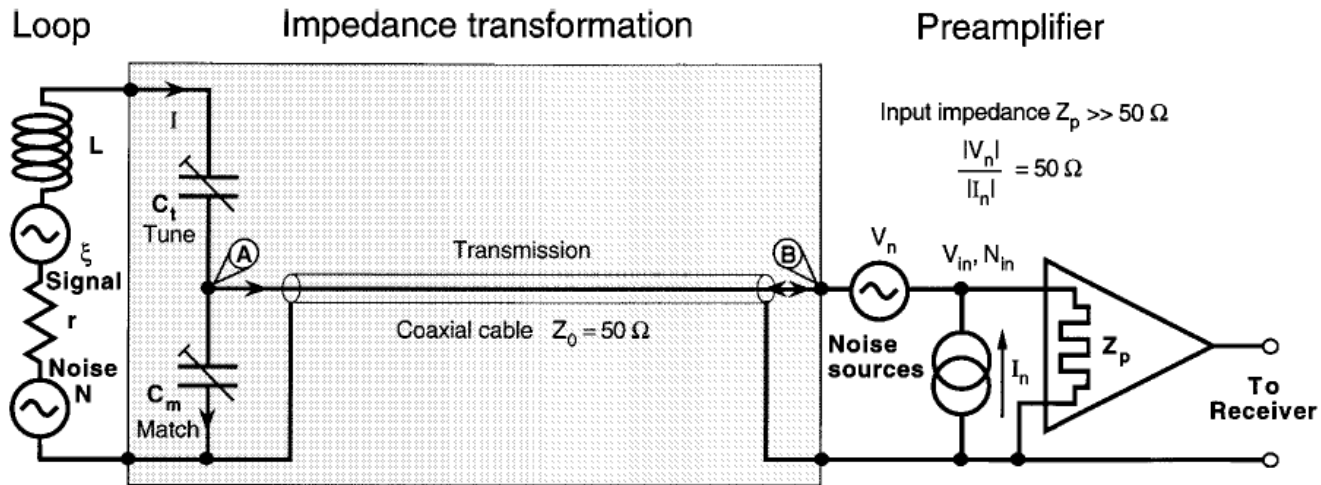
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N. J. Lambert *et al.*, Adv. Quantum Technol. 3, 1900077 (2020).

動機

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信号対雑音比



$$\Phi = \frac{\xi}{\sqrt{2rk_B(T_c + T_n)\Delta\nu}}$$

電圧増幅器の雑音温度: $T_n = \sqrt{S_{VV}(\omega)S_{II}(\omega)} \sim 75\text{K}$

標準量子限界

電圧増幅:

$$\sqrt{S_{VV}(\omega)S_{II}(\omega)} \geq \frac{\hbar\omega}{2}$$

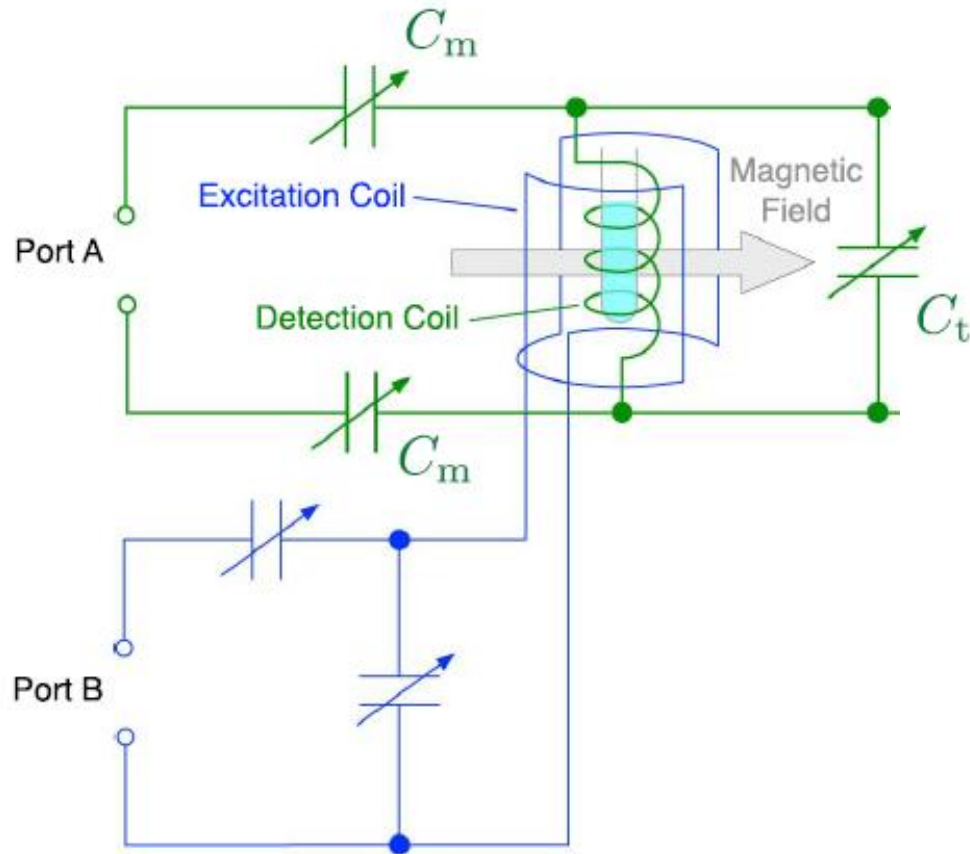
極低温では達成可

光学的位置測定:

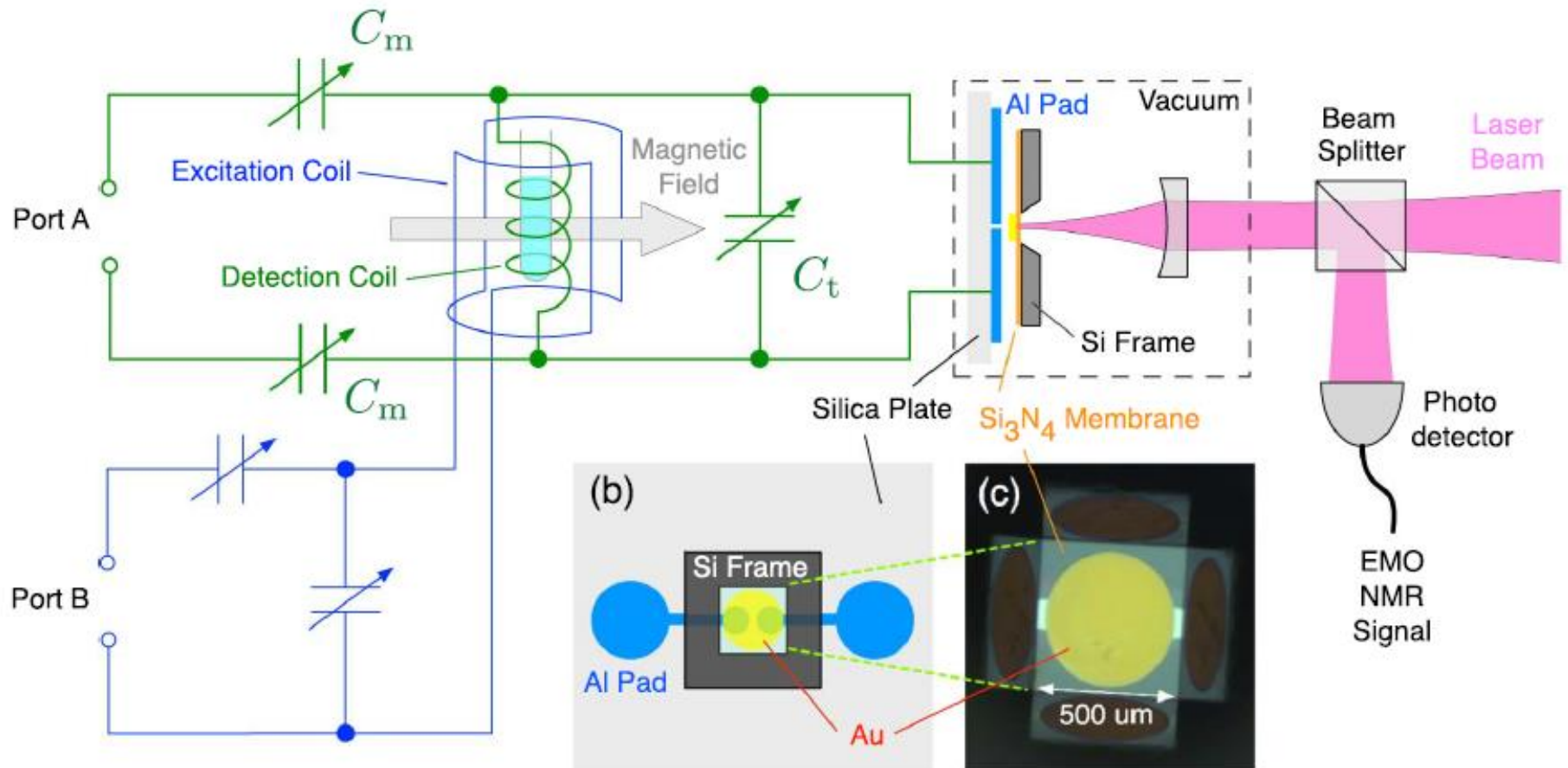
$$\sqrt{S_{FF}(\omega)S_{xx}(\omega)} \geq \frac{\hbar}{2}$$

精密光学測定で達成可

NMR信号の光検出



NMR信号の光検出

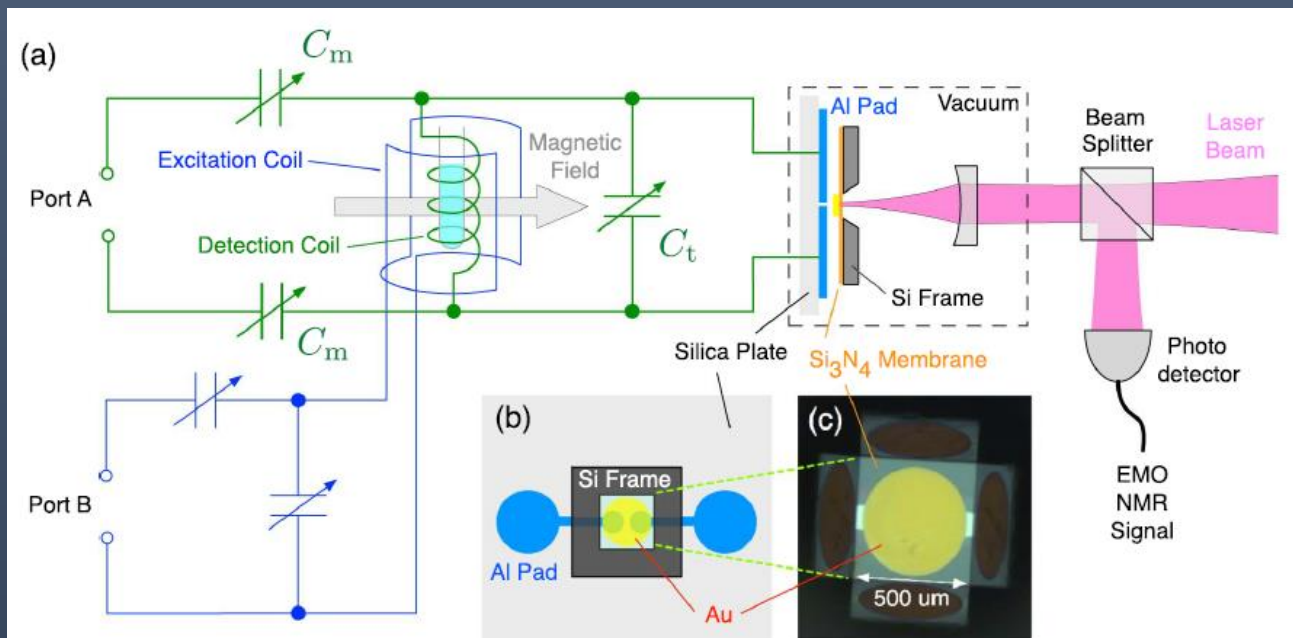
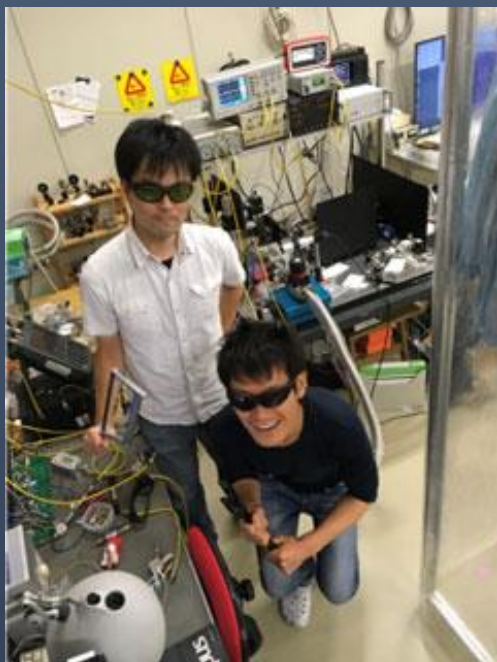


K. Takade *et al.*, Optica 5, 152 (2018)

Electro-mechano-optical detection of nuclear magnetic resonance

KAZUYUKI TAKEDA,^{1,*} KENTARO NAGASAKA,² ATSUSHI NOGUCHI,² REKISHU YAMAZAKI,²
YASUNOBU NAKAMURA,^{2,3} EIJI IWASE,⁴ JACOB M. TAYLOR,^{2,5,6} AND KOJI USAMI²

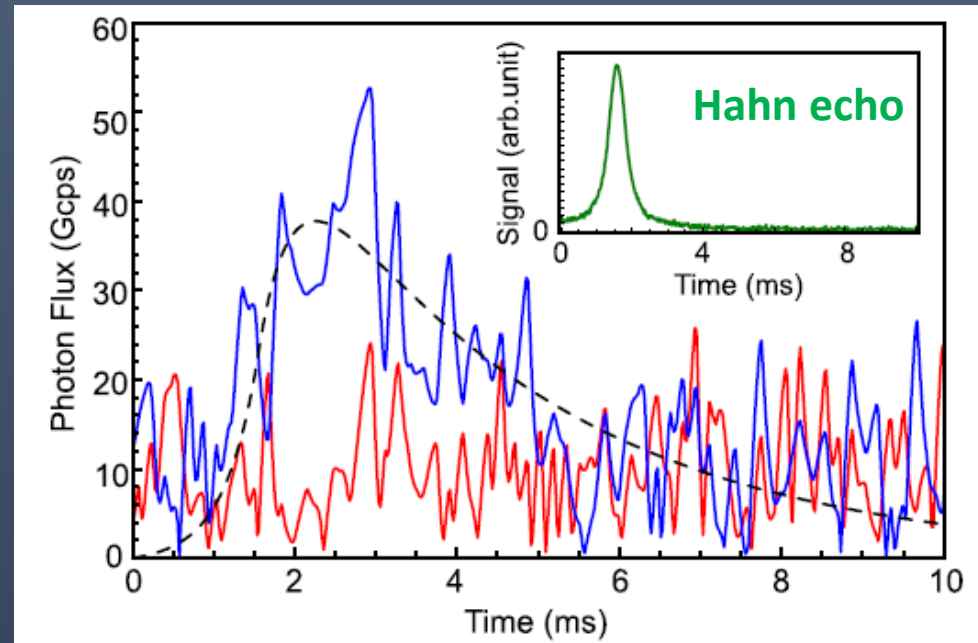
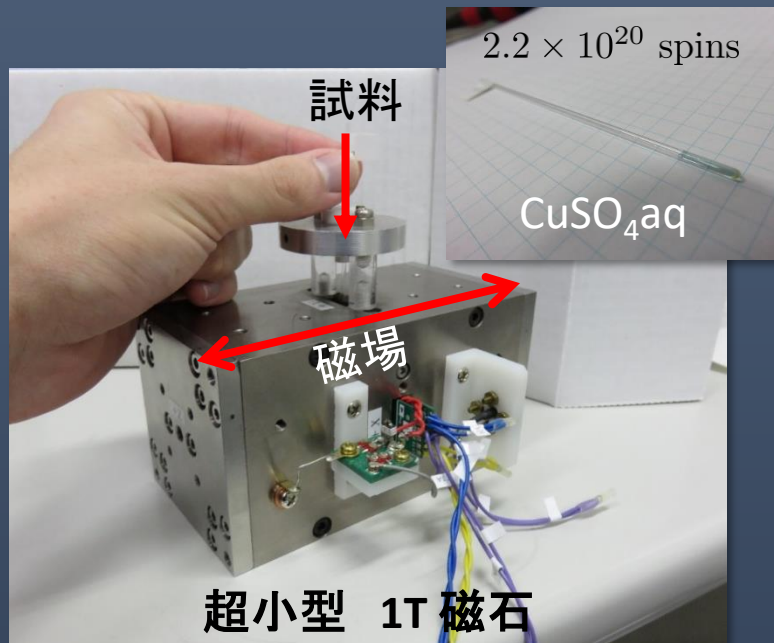
京大・化学 武田和行さん



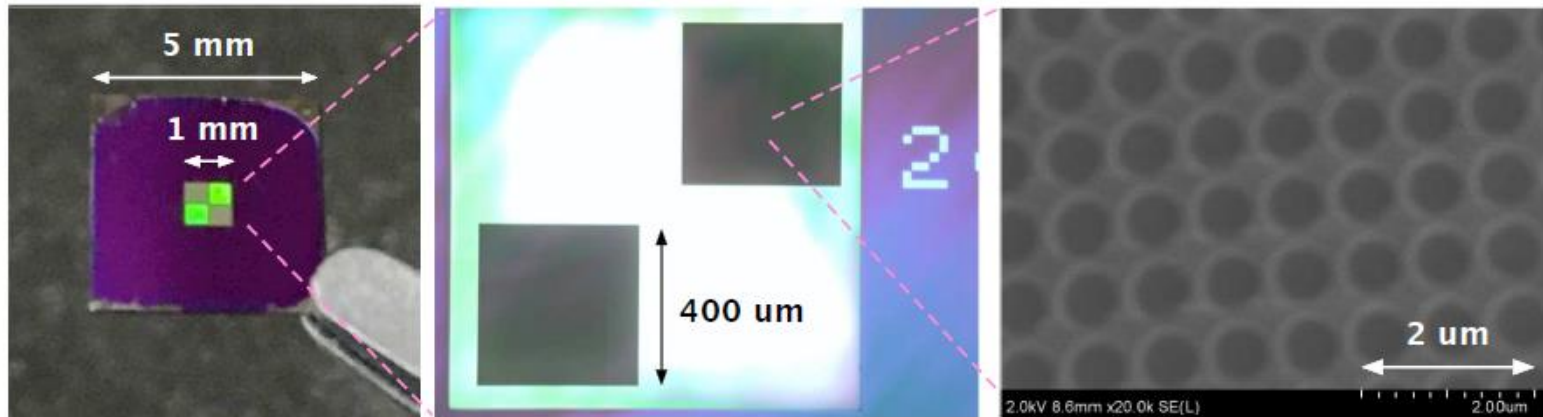
K. Takade *et al.*, Optica 5, 152 (2018)

Electro-mechano-optical detection of nuclear magnetic resonance

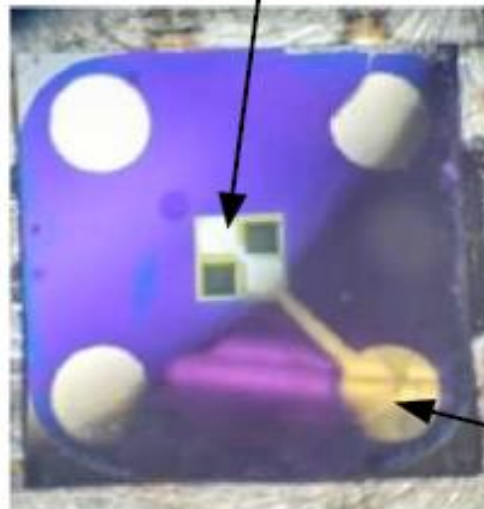
KAZUYUKI TAKEDA,^{1,*}  KENTARO NAGASAKA,² ATSUSHI NOGUCHI,² REKISHU YAMAZAKI,²
YASUNOBU NAKAMURA,^{2,3} EIJI IWASE,⁴ JACOB M. TAYLOR,^{2,5,6} AND KOJI USAMI²



メタサーフェスマイラーの採用



electrode

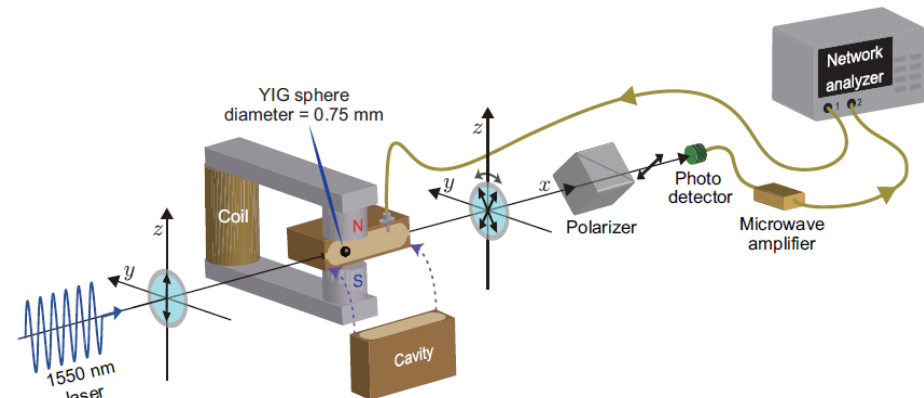


pad

量子トランスデューサー

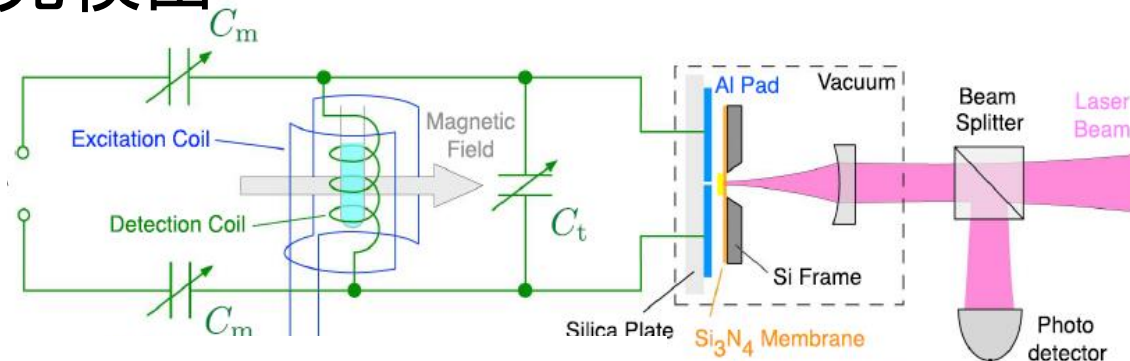
- 量子ビットのネットワーク化

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R. Hisatomi *et al.*, Phys. Rev. B **93**, 174427 (2016)

- 微弱ラジオ波信号の光検出



K. Takade *et al.*, Optica **5**, 152 (2018)